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Moment Formulas for Ensembles of Classical Compact Groups

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Introduction

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Fundamental Problem: Spacing Between Events

General Formulation: Studying system, observe values at t_1, t_2, t_3, \ldots

Question: What rules govern the spacings between the t_i ?

Examples:

- Spacings b/w Energy Levels of Nuclei.
- Spacings b/w Eigenvalues of Matrices.
- Spacings b/w Primes.
- Spacings b/w $n^k \alpha \mod 1$.
- Spacings b/w Zeros of *L*-functions.

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Classical Random Matrix Theory

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Origins of Random Matrix Theory

Classical Mechanics: 3 Body Problem Intractable.

Heavy nuclei (Uranium: 200+ protons / neutrons) worse!

Get some info by shooting high-energy neutrons into nucleus, see what comes out.

Fundamental Equation:

$$H\psi_n = E_n\psi_n$$

- H : matrix, entries depend on system
- E_n : energy levels
- ψ_n : energy eigenfunctions



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Origins of Random Matrix Theory



- Statistical Mechanics: for each configuration, calculate quantity (say pressure).
- Average over all configurations most configurations close to system average.
- Nuclear physics: choose matrix at random, calculate eigenvalues, average over matrices (real Symmetric A = A^T, complex Hermitian A^T = A).

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Random Matrix Ensembles

$$A = \begin{pmatrix} a_{11} & a_{12} & a_{13} & \cdots & a_{1N} \\ a_{12} & a_{22} & a_{23} & \cdots & a_{2N} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ a_{1N} & a_{2N} & a_{3N} & \cdots & a_{NN} \end{pmatrix} = A^{T}, \quad a_{ij} = a_{ji}$$

Fix p, define

$$\mathsf{Prob}(A) = \prod_{1 \leq i \leq j \leq N} p(a_{ij}).$$

This means

$$\operatorname{Prob}\left(\boldsymbol{A}:\boldsymbol{a}_{ij}\in[\alpha_{ij},\beta_{ij}]\right) = \prod_{1\leq i\leq j\leq N}\int_{\boldsymbol{x}_{ij}=\alpha_{ij}}^{\beta_{ij}}\boldsymbol{p}(\boldsymbol{x}_{ij})d\boldsymbol{x}_{ij}.$$

Want to understand eigenvalues of A.

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Figen	value Distribut	ion			

$$\delta(\mathbf{x} - \mathbf{x}_0)$$
 is a unit point mass at \mathbf{x}_0 :
 $\int f(\mathbf{x})\delta(\mathbf{x} - \mathbf{x}_0)d\mathbf{x} = f(\mathbf{x}_0).$

To each A, attach a probability measure:

$$\mu_{A,N}(\mathbf{x}) = \frac{1}{N} \sum_{i=1}^{N} \delta\left(\mathbf{x} - \frac{\lambda_i(A)}{2\sqrt{N}}\right)$$
$$\int_{\mathbf{a}}^{b} \mu_{A,N}(\mathbf{x}) d\mathbf{x} = \frac{\#\left\{\lambda_i : \frac{\lambda_i(A)}{2\sqrt{N}} \in [\mathbf{a}, b]\right\}}{N}$$
$$\mathbf{k}^{\text{th}} \text{ moment} = \frac{\sum_{i=1}^{N} \lambda_i(A)^k}{2^k N^{\frac{k}{2}+1}} = \frac{\text{Trace}(A^k)}{2^k N^{\frac{k}{2}+1}}.$$

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L-functions

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Riemann Zeta Function

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s} = \prod_{p \text{ prime}} \left(1 - \frac{1}{p^s}\right)^{-1}, \quad \text{Re}(s) > 1.$$

Functional Equation:

$$\xi(\mathbf{s}) = \Gamma\left(\frac{\mathbf{s}}{2}\right)\pi^{-\frac{\mathbf{s}}{2}}\zeta(\mathbf{s}) = \xi(\mathbf{1}-\mathbf{s}).$$

Riemann Hypothesis (RH):

All non-trivial zeros have
$$\operatorname{Re}(s) = \frac{1}{2}$$
; can write zeros as $\frac{1}{2} + i\gamma$.

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Genera	l/-functions				

$$L(s, f) = \sum_{n=1}^{\infty} \frac{a_f(n)}{n^s} = \prod_{p \text{ prime}} L_p(s, f)^{-1}, \quad \text{Re}(s) > 1.$$

Functional Equation:

$$\Lambda(s, f) = \Lambda_{\infty}(s, f)L(s, f) = \Lambda(1 - s, f).$$

Generalized Riemann Hypothesis (GRH):

All non-trivial zeros have
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Prope	rties of zeros	of <i>L</i> -functior	IS		

- infinitude of primes, primes in arithmetic progression.
- Chebyshev's bias: $\pi_{3,4}(x) \ge \pi_{1,4}(x)$ 'most' of the time.
- Birch and Swinnerton-Dyer conjecture.
- Goldfeld, Gross-Zagier: bound for *h*(*D*) from *L*-functions with many central point zeros.
- Even better estimates for h(D) if a positive percentage of zeros of ζ(s) are at most 1/2 − ε of the average spacing to the next zero.

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Katz-Sarnak Density Conjectures

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Measures of Spacings: *n*-Level Density and Families

Let g_i be even Schwartz functions whose Fourier Transform is compactly supported, L(s, f) an *L*-function with zeros $\frac{1}{2} + i\gamma_f$ and conductor Q_f :

$$D_{n,f}(g) = \sum_{\substack{j_1,\ldots,j_n\\j_l\neq\pm j_k}} g_1\left(\gamma_{f,j_1}\frac{\log Q_f}{2\pi}\right)\cdots g_n\left(\gamma_{f,j_n}\frac{\log Q_f}{2\pi}\right)$$

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Measures of Spacings: *n*-Level Density and Families

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Properties of *n*-level density:
 Individual zeros contribute in limit
 Most of contribution is from low zeros
 Average over similar *L*-functions (family)

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n-Leve	l Density				

n-level density: $\mathcal{F} = \bigcup \mathcal{F}_N$ a family of *L*-functions ordered by conductors, ϕ_k an even Schwartz function: $D_{n,\mathcal{F}}(\phi) =$

$$\lim_{N\to\infty}\frac{1}{|\mathcal{F}_N|}\sum_{f\in\mathcal{F}_N}\sum_{\substack{j_1,\ldots,j_n\\j_l\neq\pm j_k}}\phi_1\left(\frac{\log Q_f}{2\pi}\gamma_{j_1;f}\right)\cdots\phi_n\left(\frac{\log Q_f}{2\pi}\gamma_{j_n;f}\right)$$

As $N \to \infty$, *n*-level density converges to

$$\int \phi(\overrightarrow{\mathbf{X}})\rho_{n,\mathcal{G}(\mathcal{F})}(\overrightarrow{\mathbf{X}})d\overrightarrow{\mathbf{X}} = \int \widehat{\phi}(\overrightarrow{u})\widehat{\rho}_{n,\mathcal{G}(\mathcal{F})}(\overrightarrow{u})d\overrightarrow{u}.$$

Conjecture (Katz-Sarnak)

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(In the limit) Scaled distribution of zeros near central point agrees with scaled distribution of eigenvalues near 1 of a classical compact group.

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Corresp	ondences				

Similarities between *L*-Functions and Nuclei:

Zeros \longleftrightarrow Energy Levels

Schwartz test function \longrightarrow Neutron

Support of test function \leftrightarrow Neutron Energy.

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Theory: New Combinatorial Vantage

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Zeros	of ((s) vs GUI	=			



70 million spacings b/w adjacent zeros of $\zeta(s)$, starting at the 10^{20th} zero (from Odlyzko) versus RMT prediction.

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Key Kloosterman-Bessel integral (cuspidal newforms)

Ramanujan sum:

$${\mathcal R}(n,q) \ = \ \sum_{a ext{ mod } q}^{st} {\mathbf e}(an/q) \ = \ \sum_{d \mid (n,q)} \mu(q/d) d,$$

where * restricts the summation to be over all *a* relatively prime to *q*.

Theorem (ILS)

Let Ψ be an even Schwartz function with $\mathsf{supp}(\widehat{\Psi}) \subset (-2,2).$ Then

$$\begin{split} \sum_{n \leq N^{\epsilon}} \frac{1}{m^2} \sum_{(b,N)=1} \frac{R(m^2, b)R(1, b)}{\varphi(b)} \int_{y=0}^{\infty} J_{k-1}(y) \widehat{\Psi} \left(\frac{2\log(by\sqrt{N}/4\pi m)}{\log R} \right) \frac{\mathrm{d}y}{\log R} \\ &= -\frac{1}{2} \left[\int_{-\infty}^{\infty} \Psi(x) \frac{\sin 2\pi x}{2\pi x} \mathrm{d}x - \frac{1}{2} \Psi(0) \right] + O\left(\frac{k \log\log kN}{\log kN} \right), \end{split}$$

where $R = k^2 N$ and φ is Euler's totient function.

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$$\int_{x_1=2}^{R^{\sigma}} \int_{x_2=2}^{R^{\sigma}} \widehat{\phi}\left(\frac{\log x_1}{\log R}\right) \widehat{\phi}\left(\frac{\log x_2}{\log R}\right) J_{k-1}\left(\frac{4\pi\sqrt{m^2 x_1 x_2 N}}{c}\right) \frac{dx_1 dx_2}{\sqrt{x_1 x_2}}$$

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$$\int_{x_1=2}^{R^{\sigma}} \int_{x_2=2}^{R^{\sigma}} \widehat{\phi}\left(\frac{\log x_1}{\log R}\right) \widehat{\phi}\left(\frac{\log x_2}{\log R}\right) J_{k-1}\left(\frac{4\pi\sqrt{m^2 x_1 x_2 N}}{c}\right) \frac{dx_1 dx_2}{\sqrt{x_1 x_2}}$$

Change of variables and Jacobean:

$$\begin{array}{ccccc} u_{2} & = & x_{1}x_{2} & x_{2} & = & \frac{u_{2}}{u_{1}} \\ u_{1} & = & x_{1} & x_{1} & = & u_{1} \\ \left| \frac{\partial x}{\partial u} \right| & = & \left| \begin{array}{c} 1 & 0 \\ -\frac{u_{2}}{u_{1}^{2}} & \frac{1}{u_{1}} \end{array} \right| & = & \frac{1}{u_{1}}. \end{array}$$

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$$\int_{x_1=2}^{R^{\sigma}} \int_{x_2=2}^{R^{\sigma}} \widehat{\phi}\left(\frac{\log x_1}{\log R}\right) \widehat{\phi}\left(\frac{\log x_2}{\log R}\right) J_{k-1}\left(\frac{4\pi\sqrt{m^2 x_1 x_2 N}}{c}\right) \frac{dx_1 dx_2}{\sqrt{x_1 x_2}}$$

..

Change of variables and Jacobean:

$$\begin{aligned} u_2 &= x_1 x_2 & x_2 &= \frac{u_2}{u_1} \\ u_1 &= x_1 & x_1 &= u_1 \\ \left| \frac{\partial x}{\partial u} \right| &= \left| \begin{array}{c} 1 & 0 \\ -\frac{u_2}{u_1^2} & \frac{1}{u_1} \end{array} \right| &= \frac{1}{u_1}. \\ \int \widehat{\phi} \left(\frac{\log u_1}{\log R} \right) \widehat{\phi} \left(\frac{\log \left(\frac{u_2}{u_1} \right)}{\log R} \right) J_{k-1} \left(\frac{4\pi \sqrt{m^2 u_2 N}}{c} \right) \frac{du_1 du_2}{u_1 \sqrt{u_2}} \end{aligned}$$

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Changing variables, u_1 -integral is

$$\int_{w_1=\frac{\log u_2}{\log R}-\sigma}^{\sigma} \widehat{\phi}(w_1) \widehat{\phi}\left(\frac{\log u_2}{\log R}-w_1\right) dw_1.$$

Support conditions imply

$$\Psi_2\left(\frac{\log u_2}{\log R}\right) = \int_{w_1=-\infty}^{\infty} \widehat{\phi}(w_1) \,\widehat{\phi}\left(\frac{\log u_2}{\log R} - w_1\right) \, dw_1.$$

Substituting gives

$$\int_{u_2=0}^{\infty} J_{k-1}\left(\frac{4\pi\sqrt{m^2 u_2 N}}{c}\right) \Psi_2\left(\frac{\log u_2}{\log R}\right) \frac{du_2}{\sqrt{u_2}}$$

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n-Level Density: Katz-Sarnak Determinant Expansions

• U(N), U_k(N): det
$$(K_0(x_j, x_k))_{1 \le j,k \le n}$$

• USp(N): det
$$(\mathcal{K}_{-1}(\mathbf{x}_j, \mathbf{x}_k))_{1 \le j,k \le n}$$

• SO(even): det
$$(K_1(x_j, x_k))_{1 \le j,k \le n}$$

• SO(odd): det
$$(\mathcal{K}_{-1}(\mathbf{x}_j, \mathbf{x}_k))_{1 \le j,k \le n} + \sum_{\nu=1}^n \delta(\mathbf{x}_{\nu}) det \left(\mathcal{K}_{-1}(\mathbf{x}_j, \mathbf{x}_k)\right)_{1 \le j,k \ne \nu \le n}$$

where

$$\mathcal{K}_{\epsilon}(\mathbf{x},\mathbf{y}) = rac{\sin\left(\pi(\mathbf{x}-\mathbf{y})
ight)}{\pi(\mathbf{x}-\mathbf{y})} + \epsilon rac{\sin\left(\pi(\mathbf{x}+\mathbf{y})
ight)}{\pi(\mathbf{x}+\mathbf{y})}.$$



Manipulating determinant expansions leads to analysis of

$$\mathcal{K}(y_1,\ldots,y_n) = \sum_{m=1}^n \sum_{\substack{\lambda_1+\ldots+\lambda_m=n\\\lambda_j\geq 1}} \frac{(-1)^{m+1}}{m} \frac{n!}{\lambda_1!\cdots\lambda_m!}$$
$$\sum_{\epsilon_1,\ldots,\epsilon_n=\pm 1}^m \prod_{\ell=1}^m \chi_{\left\{ \left|\sum_{j=1}^n \eta(\ell,j)\epsilon_j y_j\right|\leq 1\right\}},$$

where

$$\eta(\ell, j) = \begin{cases} +1 & \text{if } j \leq \sum_{k=1}^{\ell} \lambda_k \\ -1 & \text{if } j > \sum_{k=1}^{\ell} \lambda_k. \end{cases}$$

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If supp
$$(\phi) \subseteq [-\frac{1}{n}, \frac{1}{n}]$$
, then each $|y_j| \leq 1/n$ and
$$\prod_{\ell=1}^m \chi_{\{|\sum_{j=1}^n \eta(\ell, j) \in_j y_j| \leq 1\}}$$

is always 1. Therefore the sum is

$$2^{n} \cdot \sum_{m=1}^{n} \sum_{\substack{\lambda_{1}+\cdots+\lambda_{m}=n\\\lambda_{j}\geq 1}} \frac{(-1)^{m+1}}{m} \frac{n!}{\lambda_{1}!\cdots\lambda_{m}!} = 0.$$

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New Results: Hughes-Miller (Orthogonal), lyer-Miller (Unitary, Symplectic): $supp(\widehat{\phi}) \subset (-\frac{1}{n-1}, \frac{1}{n-1})$

If supp
$$(\widehat{\phi}) \subseteq [-\frac{1}{n-1}, \frac{1}{n-1}]$$
, then $\left|\sum_{j=1}^{n} \eta(\ell, j) \epsilon_j y_j\right| > 1$ only when all $\eta(\ell, j) \epsilon_j y_j$ have the same sign.

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For fixed λ_j 's, m, ℓ , exactly two choices of ϵ_j 's make the product zero. Total of 2m choices as we let ℓ vary.

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For fixed λ_j 's, m, ℓ , exactly two choices of ϵ_j 's make the product zero. Total of 2m choices as we let ℓ vary.

Contributes

$$\sum_{m=1}^{n} \sum_{\substack{\lambda_1 + \dots + \lambda_m = n \\ \lambda_j \ge 1}} \frac{(-1)^{m+1}}{m} \frac{n!}{\lambda_1! \cdots \lambda_m!} (2^n - 2m) = 2(-1)^n.$$

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New Re	esult: lyer-Mille	er: Large sup	pport: $supp(\widehat{\phi}) \subseteq$	$\left[-\frac{1}{n-2}, \frac{1}{n-2}\right]$]

If supp $(\widehat{\phi}) \subseteq [-\frac{1}{n-2}, \frac{1}{n-2}]$, two new complications:

- η(ℓ, j)ϵ_jy_j need not have same sign (at most one can differ);
- more than one term in product can be zero (for fixed m, λ_j, ε_j).

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New R	esult: Iyer-Mill	er: Large s	upport: $supp(\widehat{\phi})$	$\subseteq \left[-\frac{1}{n-2}, \frac{1}{n-2}\right]$;]

If supp $(\widehat{\phi}) \subseteq [-\frac{1}{n-2}, \frac{1}{n-2}]$, two new complications:

- η(ℓ, j)ϵ_jy_j need not have same sign (at most one can differ);
- more than one term in product can be zero (for fixed m, λ_j, ε_j).

Solution: Double count (gives a contribution of $4(-1)^n$), and subtract a correcting term ρ_j .

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New Result: Iyer-Miller: Large support: $supp(\widehat{\phi}) \subseteq [-\frac{1}{n-2}, \frac{1}{n-2}]$							

Solution: Work with cumulants instead of moments:

Solution: Work with cumulants instead of moments:

$$\begin{split} C_n^{\rm SO}(\phi) &= 4(-1)^{n-1} \left(\int_{-\infty}^{\infty} \phi(x)^n \cdot \frac{\sin(2\pi x)}{2\pi x} dx - \frac{1}{2} \phi(0)^n \right) \\ &+ \frac{1}{4} \left(2n(-1)^n + \sum_{j=1}^n \rho_j \right) \\ &\times \left(\phi^n(0) - \frac{1}{2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi^{n-1}(x_1) \hat{\phi}(x_2) (1 + |x_2|) \right. \\ &\times e^{2\pi i x_1 |x_2|} \frac{\sin(2\pi x_1)}{2\pi x_1} dx_1 dx_2 \right). \end{split}$$

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Conclusion

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Conclu	ision				

- Difficulty in comparison with classical RMT is that instead of having an *n*-dimensional integral of $\phi_1(x_1) \cdots \phi_n(x_n)$ we have a 1-dimensional integral of a new test function. This leads to harder combinatorics but allows us to appeal to the result from ILS.
- Solve combinatorics by using cumulants; support restrictions translate to which terms can contribute.

