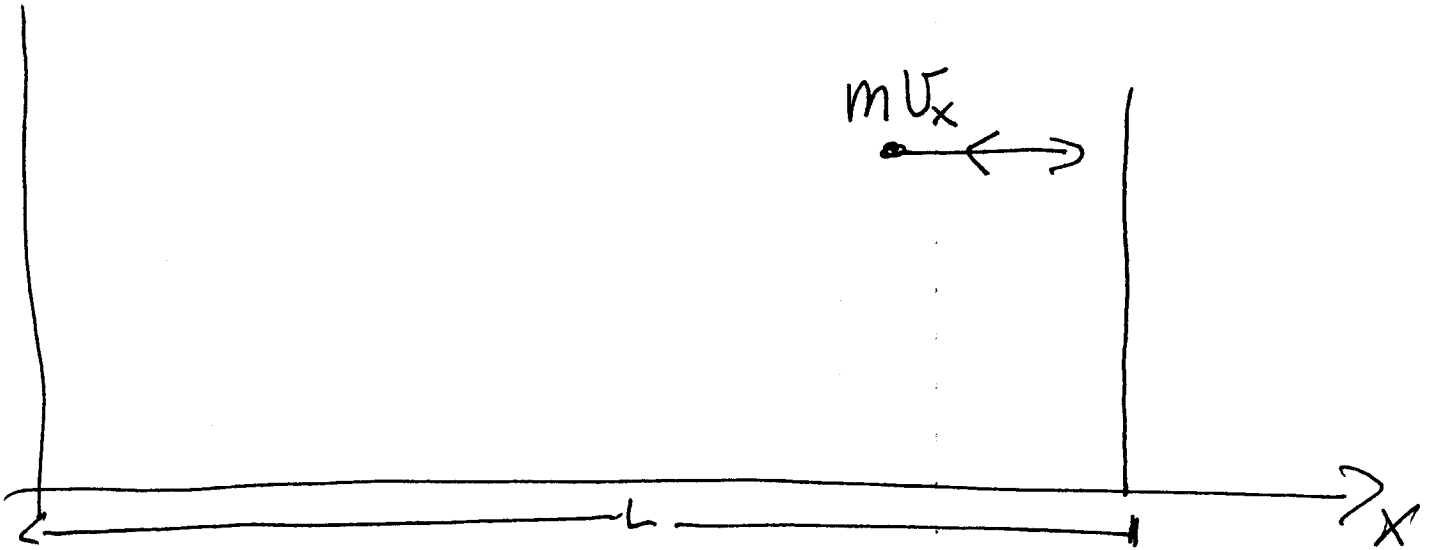


PRESSURE

$$P = \frac{F}{A} = \frac{\Delta P}{\Delta t} \frac{1}{A} = \frac{m \Delta U}{\Delta t} \frac{1}{A}$$



$$\Delta P_{x \text{ mol}} = -2mU_x$$

$$\Delta P_{x \text{ Box}} = 2mU_x$$

AFTER THE COLLISION IT WILL TAKE $\frac{2L}{U_x}$ SECONDS
 IN THIS TIME ALL THE MOL IN THE BOX
 COLLIDE ONCE WITH THE WALL
 FOR THE NEXT COLLISION WITH THE SAME
 WALL

MOMENTUM TRANSFERRED ^{TO THE WALL} PER SECOND

$$\frac{\Delta P_{\text{WALL}}}{\Delta t} = \frac{2mU_x}{\frac{2L}{U_x}} = \frac{mU_x^2}{L} \equiv F$$

For N mol

$$F = \frac{m}{L} U_{x1}^2 + \frac{m}{L} U_{x2}^2 + \dots + \frac{m}{L} U_{xN}^2$$

$$F = \frac{m}{L} N \frac{\sum_{j=1}^N U_{xj}^2}{N} = N \frac{m}{L} \overline{U_x^2}$$

AVERAGE OF THE SQUARE OF THE
X COMPONENT

$$P = \frac{F}{A} = \frac{N m \overline{U_x^2}}{V}$$

$$PV = N m \overline{U_x^2}$$

NO PREFERRED DIRECTION (SYMMETRY)

$$\overline{U_x^2} = \overline{U_y^2} = \overline{U_z^2}$$

OR

$$\overline{U^2} = 3 \overline{U_x^2}$$

$$PV = \frac{1}{3} N m \overline{U^2}$$

$$nRT = \frac{1}{3} N m \overline{u^2}$$

$$\frac{N}{N_0} RT = \frac{1}{3} N m \overline{u^2}$$

$$\frac{3RT}{N_0 M} = \overline{u^2}$$

ROOT-MEAN-SQUARE SPEED u_{rms}

$$u_{\text{rms}} = \sqrt{\frac{3RT}{M}}$$

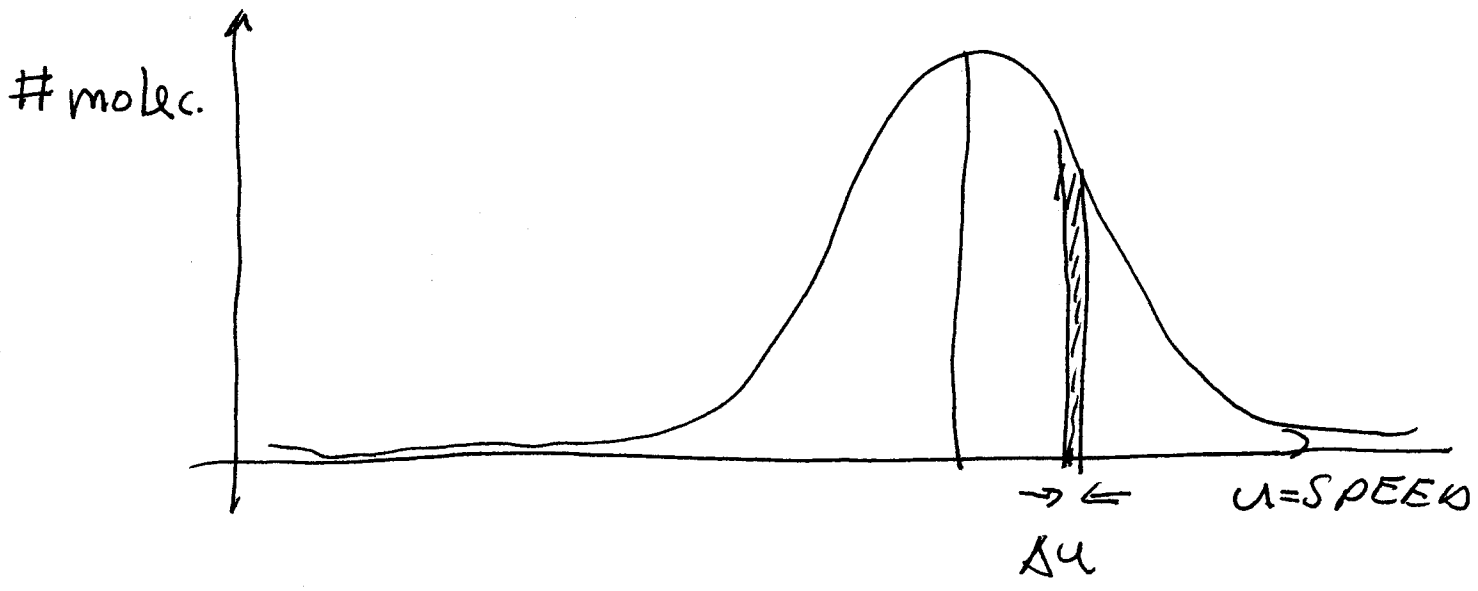
$M = N_0 M = \text{MOLAR MASS}$

AT 298 K

$$u_{\text{rms}}(\text{He}) = 1360 \frac{\text{m}}{\text{s}}$$

$$u_{\text{rms}}(\text{Xe}) = 238 \frac{\text{m}}{\text{s}}$$

DISTRIBUTION OF MOLECULAR SPEEDS



$$\frac{\Delta N}{N} = f(u) \Delta u$$

FRACTION OF MOLECULES $\Delta N/N$ THAT HAVE SPEED BETWEEN $u + \Delta u$

MAXWELL-BOLTZMANN SPEED DISTRIBUTION

$$f(u) = 4\pi \left[\frac{m}{2\pi k_B T} \right]^{3/2} u^2 e^{-\frac{mu^2}{k_B T}}$$

$k_B =$ Boltzmann const.

$$k_B = R/N_0$$

$f(u)$ IS A PROBABILITY DISTRIBUTION

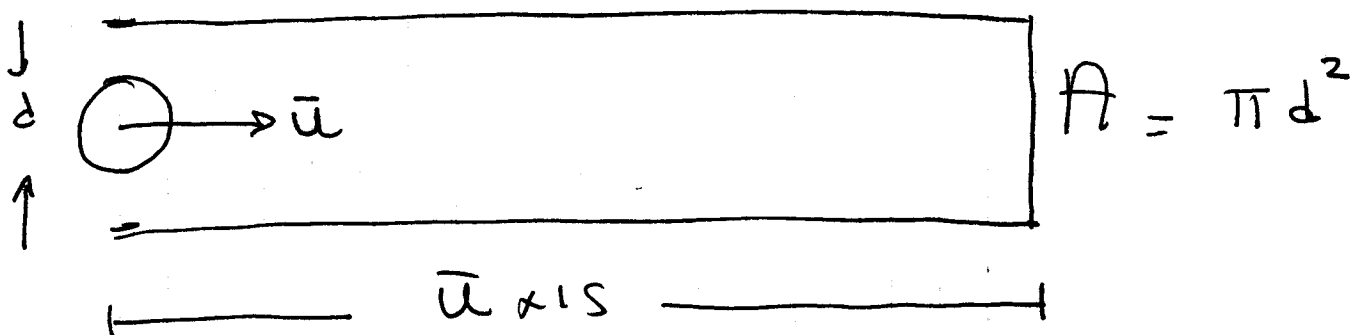
MORE PROBABLE SPEED \leftrightarrow MAX OF $f(u)$

$$u_{mp} = \sqrt{\frac{2RT}{M}}$$

AVERAGE SPEED

$$\bar{u} = \sqrt{\frac{8RT}{\pi M}}$$

FREQUENCY OF MOLECULAR COLLISIONS



$$Z_1 \propto \pi d^2 \bar{u} \frac{N}{V}$$

$$Z_1 = 4 \frac{N}{V} d^2 \sqrt{\frac{\pi RT}{M}}$$

MAX
rate constants

collisions per second

Mean free path \sim Distance between coll.

TIME BETWEEN COLLISIONS

$$\tau = \frac{1}{Z_1}$$

$$\lambda = \bar{u} \tau = \frac{1}{\sqrt{2} \pi d^2 (N/V)}$$

DIFFUSION COEFFICIENT

$$\overline{(\Delta r)^2} = 6 D t$$

FROM KTG

$$D = \frac{3\pi}{16} \bar{v} \lambda = \frac{3}{8} \sqrt{\frac{RT}{\pi M}} \frac{1}{d^2 (\text{N/V})}$$

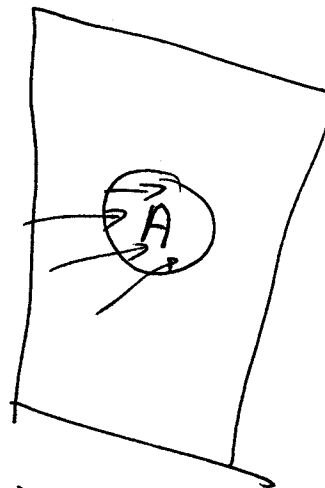
H_2 IN INTERSTELLAR SPACE

$$\lambda = 4.1 \times 10^8 \text{ m}$$

$$D = 1.4 \times 10^{11} \frac{\text{m}}{\text{s}^2}$$

OTHER RESULTSWall Collisions

$$Z_A \propto A \frac{N}{V} \bar{u}$$



EXACT CALCULATION USING $f(u)$

$$Z_A = \frac{1}{4} \frac{N}{V} \sqrt{\frac{8RT}{\pi M}} A$$

Graham's Law of Effusion

A IS A HOLE!

$$\frac{\text{rate of effusion of A}}{\text{rate of effusion of B}} = \frac{N_A}{N_B} \sqrt{\frac{M_B}{M_A}}$$

ENRICHMENT FACTOR $\sqrt{\frac{M_B}{M_A}}$ FOR THE

- LIGHTER SPECIES A

REAL GASES

$$PV = nRT \quad \text{FOR LOW PRESSURE}$$

ACTUALLY V CANNOT REACH A ZERO VALUE

$$\Rightarrow V \rightarrow V - nb$$

FOR HIGHER PRESSURES, MOLECULES SHOULD "FEEL" THE PRESENCE OF OTHER MOLECULES. THUS AN ATTRACTION OR REPULSION HAS TO OCCUR. IF WE CONSIDER AN ATTRACTIVE INTERMOLECULAR FORCE, THERE IS A TENDENCY TO LOWER THE PRESSURE, AND, TO LOWEST APPROXIMATION, THE FORCE DEPENDS ON THE PAIRING OF MOLECULES. AND THE PAIRING DEPENDS ON THE NUMBER CONCENTRATION N/V SQUARE. ($N \sim n$)

$$P = \frac{nRT}{V - nb} - a \left(\frac{n}{V} \right)^2$$

1873 J. van der Waals

$$\frac{PV}{nRT} = \frac{V}{V-nb} - \frac{a}{RT} \frac{n^2}{V}$$

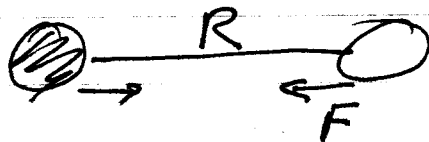
$a > 0$ ATTRACTIVE FORCE

$a < 0$ REPULSIVE FORCE

$$Z \equiv \text{COMPRESSIBILITY FACTOR} \equiv \frac{PV}{nRT}$$

INTERMOLECULAR POTENTIAL ENERGY CURVE

$V(R)$



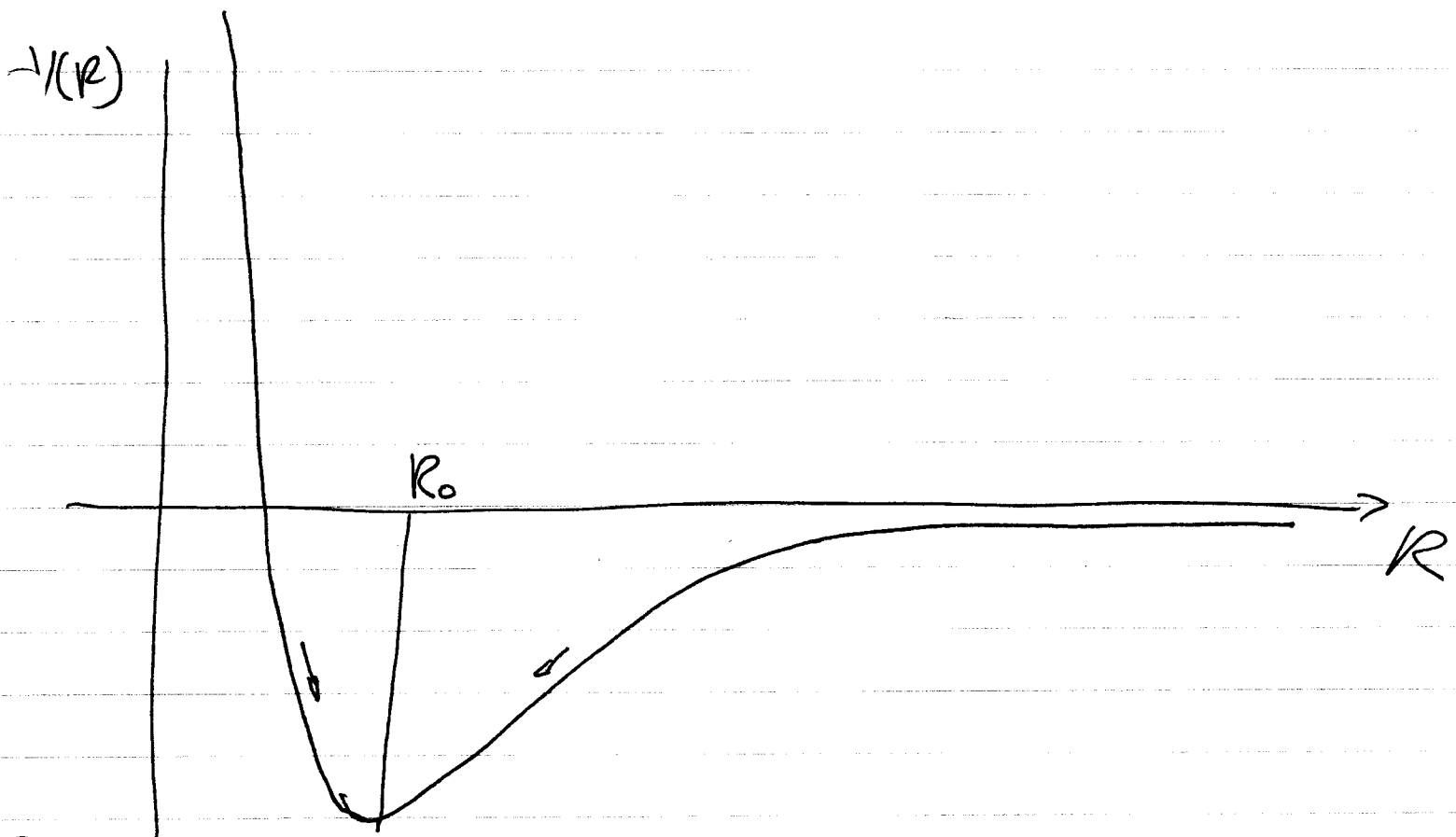
FROM PHYSICS WE KNOW THAT

$$F = - \frac{dV}{dR}$$

FORCE IS A VECTOR WHILE V IS A SCALAR FUNCTION

SHORT DISTANCE \rightarrow REPULSION

LONG DISTANCE \rightarrow ATTRACTION



R_0 "EQUILIBRIUM" DISTANCE

$$\left. \frac{dV}{dR} \right|_{R_0} = 0$$

LENNARD-JONES POT.

$$V_{LJ}(R) = 4\epsilon \left[\left(\frac{\sigma}{R} \right)^{12} - \left(\frac{\sigma}{R} \right)^6 \right]$$

CHARACTERIZATION OF GASES, LIQUIDS, AND SOLIDS.

MOLAR VOLUME

COMPRESSIBILITY

Thermal EXPANSION

Fluidity and Rigidity

Diffusion

Surface Tension

COVALENT BOND
(SHARING OF e^-)



IT TAKES 239 KJ TO BREAK A MOLE OF
Cl-Cl COVALENT BONDS

Intermolecular forces

1. - Weaker than covalent bonds
2. - less directional than covalent bonds.
3. - Longer range than covalent bonds

a) Ion-Ion \rightarrow $Na^+ Cl^-$ (Coulomb Forces)

Strong as covalent bonds

$$V(R) \sim \frac{1}{R}$$

b) Dipole-dipole $\begin{matrix} \delta^+ & & \delta^- \\ H & - & Cl \end{matrix}$

$$V(R) \sim \frac{1}{R^3}$$

ION - DIPOLE

$\text{Na}^+ \text{Cl}^-$ in H_2O

INDUCED DIPOLE FORCES

ION CAN INDUCE POLARITY TO A NONPOLAR

INDUCED DIPOLE - INDUCED DIPOLE

WEAKER

$$V(R) \sim \frac{1}{R^6}$$

BETWEEN NON POLAR DUE TO FLUCTUATIONS OF THE ELECTRONIC CHARGE = DISPERSION FORCES

REPULSIVE FORCES

Short range

$$V(R) \sim \frac{1}{R^n}$$

$$n \geq 12$$

OTHER RESULTS

LEC 3/33

SEP/14/05

DIFFUSION COEFFICIENT

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FROM KTG

$$D = \frac{3}{8} \sqrt{\frac{RT}{\pi M}} \frac{1}{d^2 (N/V)}$$

WALL COLLISIONS

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