

$$\hat{H} \Psi(x, t) = i\hbar \frac{\partial}{\partial t} \Psi(x, t)$$

$$\Psi(x, t) = a_0(t) \phi_0(x) e^{i \frac{E_0}{\hbar} t} + a_1(t) \phi_1(x) e^{i \frac{E_1}{\hbar} t}$$

$$\hat{H} = \hat{H}_0 + \hat{H}_I(t)$$

$$\hat{H}_0 \phi_i = E_i \phi_i$$

$$\hat{H}_I(t) = \frac{\mu}{2} \cdot \frac{E(t)}{2}$$

$$= \frac{\mu}{2} \cdot \frac{E_0}{2} \cos(\omega t)$$

FIRST APPROXIMATION

$$a_0(t) \approx 1$$

$$a_1(t) \approx \frac{h_{10}}{i\hbar} G_{10}(t, 0; \omega)$$

THEREFORE THE OCCUPATION PROBABILITY OF THE EXCITED STATE IS GIVEN BY:

$$P_1^{(1)}(t, \omega) = \frac{|h_{10}|^2}{\hbar^2} |G_{10}(t, 0; \omega)|^2$$

WITH

$$G_{10}(t, 0; \omega) = \int_0^t f_{\omega}(s) e^{i\omega_0 s} ds$$

WHERE

$$f_{\omega}(s) = \frac{1}{2} [e^{i\omega s} + e^{-i\omega s}] = \cos(\omega s)$$

$$\hbar\omega = E = \hbar\nu$$

$$\hbar\omega_0 = E_{10}$$

$$G_{10}(t, 0; \omega) = \frac{1}{2i} \left[\frac{e^{i(\omega_0 + \omega)t} - 1}{\omega_0 + \omega} + \frac{e^{i(\omega_0 - \omega)t} - 1}{\omega_0 - \omega} \right]$$

②

$$G_{10}(t, 0; E) = \frac{\hbar}{2i} \left[\frac{\mathcal{Q} \frac{i(E_{10}+E)t}{\hbar} - 1}{E_{10} + E} + \frac{\mathcal{Q} \frac{i(E_{10}-E)t}{\hbar} - 1}{E_{10} - E} \right]$$

$$\begin{aligned} (E_{10} + E) \frac{t}{\hbar} &= 2\pi (\cancel{1} + x) \frac{E_{10} t}{\hbar} \\ &= 2\pi (1+x) \frac{t}{(\hbar/E_{10})} = 2\pi (1+x) \tau \end{aligned}$$

$$ut \equiv \frac{\hbar}{E_{10}} = \frac{4. \text{ eV fs}}{E_{10}} \approx \frac{4 \text{ eV}}{3 \text{ eV}} \text{ fs}$$

$$ut \approx \text{fs} \quad x = E/E_{10}$$

$$G_{10}(\tau, 0; x) = \left(\frac{\hbar}{E_{10}} \right) \left(\frac{1}{4\pi i} \right) \left[\frac{\mathcal{Q} \frac{i2\pi(1+x)\tau}{\hbar} - 1}{1+x} + \frac{\mathcal{Q} \frac{i2\pi(1-x)\tau}{\hbar} - 1}{1-x} \right]$$

$$\tau = \frac{t}{ut} \sim \text{Time in fs}$$

$$x = \frac{E}{E_{10}}$$

$$\hat{H} = \hat{H}_0 + \hat{H}_I$$

$$\hat{H}_0 = \hat{H}_{\text{vib}} + \hat{H}_{\text{rot}}$$

$$E_{n\ell m} = \hbar\omega\left(\frac{1}{2} + n\right) + \hbar^2 \ell(\ell+1)$$

$$\textcircled{\text{I}} \quad H_{\text{I}}^{\text{vib}} \sim \frac{E \cdot \mu}{2} \sim |E_0| \frac{\mu}{\kappa} \quad \text{VIB}$$

$$\textcircled{\text{II}} \quad H_{\text{I}}^{\text{rot}} \sim \frac{E \cdot \mu}{2} \sim |E_0| \frac{\mu}{2} \quad \text{ROT}$$

TRANSITIONS

$$\langle n' | H_{\text{I}}^{\text{vib}} | n \rangle \quad \text{OR} \quad \langle \ell', m' | H_{\text{I}}^{\text{rot}} | \ell, m \rangle$$

(I) Vibrational

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$$h_{10} \equiv \int_{\Omega} \phi_1^*(x) \left(\hat{h}(x) \phi_0(x) \right) dx \equiv \langle 1 | \hat{h} | 0 \rangle$$

$$H_F(x, t) = -q \frac{\mu}{\epsilon_0} \cdot \underline{E}(t) = -q \frac{\mu}{\epsilon_0} E_0 \cos(\omega t)$$

$$\hat{h}(x) \equiv q \frac{\mu}{\epsilon_0} E_0 = \overbrace{q |E_0| \frac{\mu}{\epsilon_0}(x)} \text{ SF } E_0 = \hat{x} |E_0|$$

$$h_{10} = q |E_0| \int_{\Omega} \phi_1^*(x) \left(\frac{\mu}{\epsilon_0}(x) \phi_0(x) \right) dx \equiv q |E_0| \langle 1 | \hat{x} | 0 \rangle$$

$$\frac{\mu}{\epsilon_0}(x) = \mu^0 + \left. \frac{d\mu}{dx} \right|_0 x + \frac{1}{2} \left. \frac{d^2\mu}{dx^2} \right|_0 x^2 + \dots$$

Transitions

$$\begin{aligned} \langle n' | \hat{x} | n \rangle &= \mu^0 \langle n' | n \rangle + \left. \frac{d\mu}{dx} \right|_0 \langle n' | \hat{x} | n \rangle \\ &+ \frac{1}{2} \left(\left. \frac{d^2\mu}{dx^2} \right|_0 \right) \langle n' | \hat{x}^2 | n \rangle + \dots \end{aligned}$$

Permanent dipole $\Rightarrow \mu^0 \neq 0$

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IF $n' \neq n$ AND $\left. \frac{dU(x)}{dx} \right|_0 \neq 0$ (SR ACTIVE)
 $\langle n' | n \rangle = 0$

AND

Vib

$$\langle n' | \hat{x} | n \rangle \neq 0$$

WHERE

$$|n\rangle \sim H_n(x) e^{-\frac{\alpha x^2}{2}}$$

WE KNOW THAT

$$x H_n(x) \sim a H_{n-1}(x) + b H_{n+1}(x)$$

THEREFORE

$$\langle n' | \hat{x} | n \rangle \sim a \langle n' | n-1 \rangle + b \langle n' | n+1 \rangle$$

$$\Rightarrow \Delta n' = \pm 1$$

OR

$$n \rightarrow n+1$$

$$n \rightarrow n-1$$

ANYTHING ELSE IS FORBIDDEN IN THE
HO APPROXIMATION