

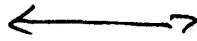
DIFFERENT
CONSTRAINTS

N, V, E

N, V, T

μ, V, T

N, P, T



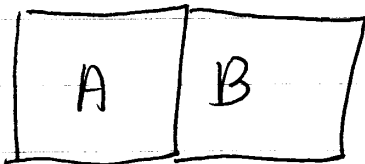
DIFFERENT
ENSEMBLE
CANONICAL
CANONICAL
(VARIABLE E_T)

GRAND CANONICAL
(VARIABLES E AND N)

PRESSURE ENSEMBLE
(VARIABLE V AND E)

NON INTERACTING SYSTEMS

$$Q = \sum_n \Omega e^{-\beta E_n}$$



FOR TWO SYSTEMS

$$E_{ABmn} \approx E_{An} + E_{Bm} + \int_{ABMN}$$

FOR IDEAL GAS-TYPE SYSTEMS

$$E_{ABmn} \approx E_{An} + E_{Bm}$$

$$Q = \sum_m \sum_n \Omega e^{-\beta E_{An} - \beta E_{Bm}} = q_A q_B$$

$$q_A = \sum_n \Omega e^{-\beta E_{An}}$$

$$Q = q_A q_B \quad \text{IF } A=B \quad Q = q^2$$

FOR N SYSTEMS

$$Q = q^N$$

DISTINGUISHABLE

FOR N PARTICLES

$$Q = \frac{1}{N!} q^N$$

INDISTINGUISHABLE

NON INTERACTING (IDEAL) PARTICLES.
(WEAK APPROXIMATION)

FOR N molecules

$$Q = \frac{1}{N!} q_{\text{mol}}^N$$

$$E_{\text{mol}} = E_{\text{TRA}} + E_{\text{ROT}} + E_{\text{VIB}} + E_{\text{elec.}}$$

$$q_{\text{mol}} = q_T q_R q_V q_E$$

THIS IS BETTER APPROX. SINCE THERE IS
NO INTERFERENCE BETWEEN DEGREE
OF FREEDOM

TRANSLATIONAL MOTION

QU → PARTICLE IN A BOX

$$E_n = \frac{h^2}{8mL^2} n^2$$

$$\Delta E_{n, n+1} = \frac{h^2}{2mL^2} (2n+1)$$

$$\frac{h^2}{2mL} \sim 10^{-15} \text{ cm}^{-1} \cong 10^{-38} \text{ J}$$

$$\text{For } L \sim 10^{-2} \text{ m}$$

FOR MACRO h WE HAVE A CONTINUUM OF ENERGY

$$\sum \mathcal{Q}^{-\beta E_n} \rightarrow \int_0^{\infty} \mathcal{Q}^{-\frac{\beta h^2}{8mL^2} n^2} dn$$

$$\int_0^{\infty} \mathcal{Q}^{-\alpha x^2} dx = \frac{1}{2} \sqrt{\frac{\pi}{\alpha}}$$

$$q_T = \int_0^{\infty} A e^{-\frac{\beta \hbar^2}{8mL^2} x^2} dx$$

$$= \frac{1}{2} \sqrt{\frac{\pi 8mL^2}{\beta \hbar^2}} = \frac{1}{2} \frac{L}{\hbar} \sqrt{8m k_B T \pi}$$

$$\boxed{q_T = \frac{L}{\Lambda}} \quad 1-d$$

$$\Lambda = \frac{h}{\sqrt{2\pi m k_B T}}$$

THERMAL
WAVELENGTH

FROM KINETIC THEORY

$$\frac{1}{2} \frac{\bar{p}^2}{m} = \frac{1}{2} k_B T \Rightarrow \bar{p}^2 = m k_B T$$

$$\Lambda = \frac{h}{\bar{p}}$$

in 3-d

$$q_T = \frac{V}{\Lambda^3} = (2\pi m k_B T)^{3/2} \frac{V}{h^3}$$

ROTATIONAL MOTION

$$q_R = \sum_{J=0}^{\infty} (2J+1) \mathcal{Q}^{-\beta h c B J(J+1)}$$

APPROX $\rightarrow \int_0^{\infty} (2X+1) \mathcal{Q}^{-\beta h c B X(X+1)} dx$

$$q_R = \frac{1}{h c \beta B} = \frac{k_B T}{h c B}$$

$$q_R = \frac{T}{\Theta_R}$$

$$\Theta_R = \frac{h c B}{k_B}$$

ROTATIONAL
TEMPERATURE

SYMMETRY CONSIDERATIONS

HOMONUCLEAR DIATOMIC MOLECULES
HETERONUCLEAR

$$q_r = \frac{T}{\sigma \Theta_r}$$

AT LOW TEMPERATURES THE INTEGRAL IS NOT A GOOD APPROXIMATION OF THE SUM!

POLYATOMIC MOLECULES

ALL MOLECULES HAVE 3 PRINCIPAL AXIS OF ROTATION!

$\Rightarrow I_i$

$$q_r = \frac{\sqrt{\pi}}{\sigma} \left[\frac{1}{\beta h c B_A} \right]^{1/2} \sqrt{\frac{1}{\beta h c B_B}} \sqrt{\frac{1}{\beta h c B_C}}$$

WHERE $B_i = \frac{h}{8\pi^2 I_i}$

$$k_B T \approx \Delta E_{\text{rot}} \quad \text{AT ROOM TEMP}$$

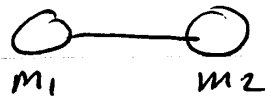
\Rightarrow DISTRIBUTION AND
POPULATION OF HIGHER
ROTATIONAL ENERGY
LEVELS

NOTE:

$$p_i = \frac{\Omega^{-\beta \epsilon_i}}{\sum_{n=0}^{\infty} \Omega^{-\beta \epsilon_n}} = \frac{\Omega^{-\beta(\epsilon_i - \epsilon_0)}}{1 + \sum_{n=1}^{\infty} \Omega^{-\beta(\epsilon_n - \epsilon_0)}}$$

THE PARTITION FUNCTION ONLY DEPENDS
ON THE ENERGY DIFFERENCES!

ROTATIONAL MOTION



DIATOMIC MOLECULE

$$E_J = B J(J+1) hc \quad J = 0, 1, 2, \dots$$

$$B = \frac{h}{8\pi^2 c I}$$

$$I = \mu R^2$$

$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$

$$g_J = 2J + 1$$

$$q_R = \sum_{J=0}^{\infty} (2J+1) e^{-\beta \cdot hc B J(J+1)}$$

$$\rightarrow \int_0^{\infty} (2x+1) e^{-\beta hc B x(x+1)} dx$$

$$\frac{d}{dx} e^{-\beta hc B x(x+1)} = -\beta hc B (2x+1) e^{-\beta hc B x(x+1)}$$

$$\boxed{q_R = \frac{1}{\beta hc B} = \frac{k_B T}{hc B}}$$

VIBRATIONAL PARTITION FUNCTION

HO - QM

$$E_n = hc\bar{\nu} \left(\frac{1}{2} + n \right) \quad n=0, 1, \dots$$

$$q_v = \sum_{n=0}^{\infty} \mathcal{Q}^{-\beta E_n} = \mathcal{Q}^{-\frac{\beta hc\bar{\nu}}{2}} \sum_{n=0}^{\infty} \mathcal{Q}^{-\beta hc\bar{\nu} n}$$

$$= \mathcal{Q}^{-\frac{\beta hc\bar{\nu}}{2}} \sum_{n=0}^{\infty} \left[\mathcal{Q}^{-\beta hc\bar{\nu}} \right]^n$$

$$\sum_{n=0}^{\infty} X^n = \frac{1}{1-X} \approx 1 + X + \dots$$

$$q_v = \frac{\mathcal{Q}^{-\frac{\beta hc\bar{\nu}}{2}}}{1 - \mathcal{Q}^{-\beta hc\bar{\nu}}} \quad \begin{array}{l} \rightarrow \frac{1}{1 - \mathcal{Q}^{-\beta hc\bar{\nu}}} \\ \downarrow \text{ENERGY SHIFT} \\ \Rightarrow E_0 = 0 \end{array}$$

$\downarrow E_0 \neq 0$