

$$\Gamma(s)\Gamma(1-s) = \frac{\pi}{\sin \pi s}$$

$$\Gamma(s)\Gamma(s+1/2) = 2^{1-2s} \pi^{1/2} \Gamma(2s)$$

$$\zeta(s) = \frac{1}{2}s(s-1)\Gamma(\frac{s}{2})\pi^{-s/2}f(s) = \overline{\zeta}(1-s)$$

$$\hookrightarrow \text{Start} + \int_0^\infty x^{\frac{s}{2}-1} e^{-n^2 \pi x} dx = \frac{\Gamma(s/2)}{n^s \pi^{s/2}}$$

$$\pi^{-\frac{s}{2}} \Gamma(\frac{s}{2}) f(s) = \int_0^\infty x^{\frac{s}{2}-1} \omega(x) dx$$

$$\omega(x) = \sum_{n=1}^{\infty} e^{-n^2 \pi x} \xrightarrow{x \rightarrow \infty} 0$$

$$\omega'(x) = -\frac{1}{2} - \frac{1}{2} x^{\frac{1}{2}} + x^{\frac{1}{2}} \omega(x)$$

$$\pi^{-\frac{s}{2}} \Gamma(\frac{s}{2}) f(s) = \frac{1}{s(s-1)} + \int_1^\infty (x^{\frac{s}{2}-1} + x^{-\frac{s}{2}-\frac{1}{2}}) \omega(x) dx$$

Simple Proof: $\operatorname{Re}(s) > 0$:

$$f(s) = \frac{s}{s-1} - s \int_1^\infty \frac{x^s}{x^{s+1}} dx$$

Method of Stationary Phase

Laplace's Method: Pg 323

Consider $\int_a^b e^{-s\Phi(x)} \cdot \psi(x) dx$ where phase $\Phi(x)$ is real and Φ, ψ infinitely differentiable (for convergence).

Suppose Φ has a minimum at $x_0 \in (a, b)$, so $\Phi'(x_0) = 0$, and also $\Phi''(x_0) > 0$ in $(a, b]$. For $s \rightarrow \infty$ have

$$\int_a^b e^{-s\Phi(x)} \psi(x) dx = e^{-s\Phi(x_0)} \left[\frac{A}{s^{1/2}} + O(\frac{1}{s}) \right]$$

$$\text{with } A = \sqrt{2\pi} \frac{\psi(x_0)}{\sqrt{\Phi''(x_0)}}$$

Step 1: Why $\Phi(x_0) = 0$ in proof (study $\Phi(x) - \Phi(x_0)$)

Step 2: Taylor: $\Phi(x) = \frac{1}{2} \Phi''(x_0) (x - x_0)^2 + \psi(x)$ with $\psi(x_0) = 1 + O(x - x_0)$
as $x \rightarrow x_0$

Step 3: Change vars: $x \mapsto y = (x - x_0) \Phi(x)^{1/2}$
 \hookrightarrow note $\frac{dy}{dx}|_{x_0} = 1$ so $\frac{dx}{dy} = 1 + O(y)$ as $y \rightarrow 0$

\hookrightarrow have $\psi(x) = \tilde{\psi}(y) = \psi(x_0) + O(y)$ as $y \rightarrow 0$

Step 4: $[a', b'] \ni x_0$ Then, with $\alpha < 0 < \beta$
 $\int_{a'}^{b'} e^{-s\Phi(x)} \psi(x) dx = \psi(x_0) \int_\alpha^\beta e^{-s\frac{\Phi''(x_0)}{2} y^2} dy + O\left(\int_\alpha^\beta e^{-s\frac{\Phi''(x_0)}{2} y^2} |y| dy\right)$

Step 5: Different than book:

$$\text{use } \int_{-\infty}^{\infty} e^{-y^2/2\sigma^2} dy = \sqrt{2\pi} \cdot \sigma$$

with $\sigma^2 = \sqrt{s\Phi''(x_0)}$, "small" error with $-\alpha = \beta = \sigma$

Method of Stationary Phase

Application: Stirling's Formula

See power & continuation of factorial to understand it at integer values. Study

$$\Gamma(n+1) = \int_0^\infty e^{-x} x^n dx = n! \quad (\text{so } s = n+1)$$

$$\Gamma(s+1) = \int_0^\infty x^s e^{-x} dx$$

$$= \int_0^\infty e^{s \log x} \cdot e^{-x} dx$$

$$= \int_0^\infty e^{-s(-\log x)} e^{-x} dx$$

$\Phi(x) = -\log x$ } doesn't
 $\psi(x) = e^{-x}$ } work

Better

$$\Gamma(s) = \int_0^\infty e^{-x} x^s \frac{dx}{x}$$

$$= \int_0^\infty e^{-x+s \log x} \frac{dx}{x} \quad \text{change var: } x \mapsto xs \text{ so } \frac{dx}{x} \mapsto \frac{dx}{x}$$

$$= \int_0^\infty e^{-sx+s \log sx} \frac{dx}{x}$$

$$= e^{s \log s} e^{-s} \int_0^\infty e^{-s\Phi(x)} \frac{dx}{x} \quad \text{with } \Phi(x) = x - 1 - \log x$$

Why this change?

$$\Phi'(1) = \Phi'(1) = 0$$

$$\Phi''(x) = \sqrt{x^2} > 0 \text{ so } \Phi''(x_0) = 1 \text{ as } x_0 = 1$$

$$\psi(x) = \sqrt{x} \text{ so } \psi(x_0) = 1 \text{ as } x_0 = 1$$

$$\text{Gauß theorem: } A = \sqrt{2\pi} \cdot \sqrt{\sqrt{1}} = \sqrt{2\pi}$$

$$\Gamma(s) \sim e^{s \log s} e^{-s} \cdot e^{-s\Phi(1)} \cdot \frac{A}{s^{1/2}}$$

$$\Gamma(s) \sim s^s e^{-s} \sqrt{\frac{2\pi}{s}} (1 + \text{small})$$

$$\Gamma(n+1) = n \Gamma(n) \approx n^n e^{-n} \sqrt{2\pi n} (1 + \text{small})$$