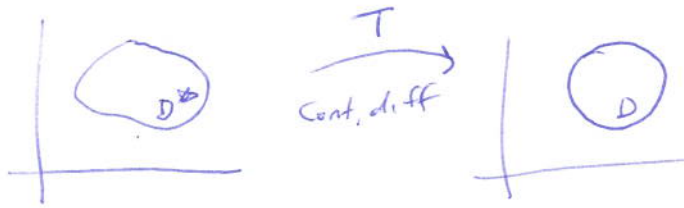


CHAPTER 6: THE CHANGE OF VARIABLES FORMULA + APPS TO $\int$

6.1. THE GEOMETRY OF MAPS FROM $\mathbb{R}^2$ TO $\mathbb{R}^2$



ex: $T: [0,1] \times [0,2\pi] \rightarrow \text{unit circle} : T(r,\theta) = (r\cos\theta, r\sin\theta)$

$\hookrightarrow$ not 1-1 (inj): $\{0\} \times [0,2\pi] \rightarrow (0,0)$

is onto (surj): all points in D hit

Thm: A 2×2 matrix, $\text{Det } A \neq 0$, maps $\square$ to $\square$, vertices $\rightarrow$ vertices

Thm (Adj): D^* , D elemental, $T: D^* \rightarrow D$ st $DT(u,v) \neq \vec{0} \ \forall (u,v) \in D^*$,
if $T: \partial D^* \xrightarrow[\text{onto}]{1-1} \partial D$ Then T is 1-1, onto (bij)

Thm: $\text{Det } A \neq 0 \iff A$ is 1-1, onto

HW: pg 375: #2, #7, #12

Suggested: #3, #8

6.2. THE CHANGE OF VARIABLE FORMULA

$T: D^* \rightarrow D$ with $T(u,v) = (x,y) = (x(u,v), y(u,v))$

Can't be $\iint_D f(x,y) dx dy = \iint_{D^*} f(x(u,v), y(u,v)) du dv$

↳ in general D, D^* diff areas, take $f \equiv 1$

Defn: Jacobian Det

$$"JF" = DT(u_0, v_0) = \begin{vmatrix} \frac{\partial x}{\partial u}(u_0, v_0) & \frac{\partial x}{\partial v}(u_0, v_0) \\ \frac{\partial y}{\partial u}(u_0, v_0) & \frac{\partial y}{\partial v}(u_0, v_0) \end{vmatrix} = \det \left| \frac{\partial(x,y)}{\partial(u,v)} \right|$$

$$\text{Yields } \iint_D f(x,y) dx dy = \iint_{D^*} f(x(u,v), y(u,v)) \left| \frac{\partial(x,y)}{\partial(u,v)} \right| du dv$$

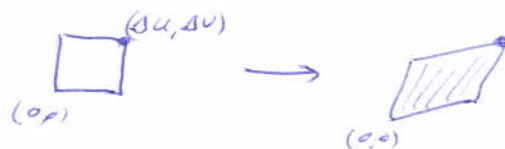
↳ Think of as $\frac{dx dy}{du dv}$
assuming T is $C^1, 1-1$

Justification

Pg 134: Differentiability $g(\vec{v}) = g(\vec{v}_0) + Dg(\vec{v}_0) \cdot (\vec{v} - \vec{v}_0) + \text{small}$

w/ $g(u_0, v_0) = (x_0, y_0) = (0, 0)$, $T(u_0, v_0) = (x_0, y_0)$

$$T(\Delta u, \Delta v) \approx \begin{pmatrix} \frac{\partial x}{\partial u} & \frac{\partial x}{\partial v} \\ \frac{\partial y}{\partial u} & \frac{\partial y}{\partial v} \end{pmatrix} \begin{pmatrix} \Delta u \\ \Delta v \end{pmatrix} + \text{small}$$

 so Area $\Delta x \Delta y \approx |\det DT| \cdot \Delta u \Delta v$

Change of vars: 1-var

$$\int_a^b f(x(u)) \cdot x'(u) du = \int_{x(a)}^{x(b)} f(x) dx \quad \text{Chain Rule}$$

6.2. The Change of Variables Formula

Polar Coords: $\iint_D f(x,y) dx dy = \iint_{D'} f(r \cos \theta, r \sin \theta) \cdot r dr d\theta$

↳ not 1-1 everywhere, but ok

Ex: Gaussian integral (and $\Gamma(1/2)$ and Wallis' Factorials...)

$$\int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{\pi} : \text{Square and change to polar}$$

Triple and higher

$$\begin{vmatrix} \partial x_1 / \partial u_1 & \dots & \partial x_1 / \partial u_n \\ \vdots & & \vdots \\ \partial x_n / \partial u_1 & \dots & \partial x_n / \partial u_n \end{vmatrix} = \left| \frac{\partial (x_1, \dots, x_n)}{\partial (u_1, \dots, u_n)} \right|$$

Cylindrical: $r dr d\theta dz$ $x = r \cos \theta, y = r \sin \theta, z = z$

Spherical: $\rho^2 \sin \phi d\rho d\phi d\theta$ $x = \rho \sin \phi \cos \theta, y = \rho \sin \phi \sin \theta, z = \rho \cos \phi$

HW: pg 390: #2, #3a, #11, #29, #34

Suggested: #4, #7, #10, #15, #18, #31

6.4. IMPROPER INTEGRALS

One variable: $\int_0^1 x^{-1/2} dx$ or $\int_1^\infty x^{-2} dx$

→ use limits: $\lim_{\epsilon \rightarrow 0} \int_\epsilon^1 x^{-1/2} dx$

Several vars: look at contraction, expand

$$D = \{(x, y): a \leq x \leq b, \phi_1(x) \leq y \leq \phi_2(x)\}$$

$$D_{\eta, \delta} = \{(x, y): a + \eta \leq x \leq b - \eta, \phi_1(x) + \delta \leq y \leq \phi_2(x) - \delta\}$$

$$\lim_{(\eta, \delta) \rightarrow (0, 0)} \iint_{D_{\eta, \delta}} f dA$$

if this exists, then $\lim_{\eta \rightarrow 0} \lim_{\delta \rightarrow 0} \int_{a+\eta}^{b-\eta} \int_{\phi_1(x)+\delta}^{\phi_2(x)-\delta} f dA$ exists

However, iterated limits exist $\nrightarrow \lim_{(\eta, \delta) \rightarrow (0, 0)}$ exists

→ example in book

Fubini: Delemersta!, $f \geq 0$, $\left[\iint_D f dA, \int_a^b \int_{\phi_1(x)}^{\phi_2(x)} f dy dx \text{ for } y\text{-simple} \right.$
or $\left. \int_c^d \int_{\psi_1(y)}^{\psi_2(y)} f dx dy \text{ for } x\text{-simple exist} \right]$ then f integrable
and all equal.

→ advanced: $\sum \sum a_{mn} \neq \sum \sum a_{nm}$

$$\begin{matrix} 0 & 0 \\ 0 & -1 \\ -1 & +1 \\ +1 & 0 & \dots \end{matrix}$$

Hw: pg 415: #5, #17, #18

Suggested: #1, #12, #13

Gamma Function: $\Gamma(n) = (n-1)!$

$$\Gamma(s) = \int_0^{\infty} x^{s-1} e^{-x} dx \quad (= 2 \int_0^{\infty} e^{-t^2} t^{2s-1} dt)$$

$$\Gamma(x)\Gamma(-x) = -\frac{\pi}{x \sin(\pi x)}$$

$$|(ix)!|^2 = \frac{\pi x}{\sinh(\pi x)}$$

$$\Gamma(2x) = (2\pi)^{-1/2} 2^{2x-1/2} \Gamma(x) \Gamma(x+\frac{1}{2})$$

$$\Gamma(\frac{1}{2}) = \sqrt{\pi}$$

$$\text{Stirling: } \Gamma(z) \sim e^{-z} z^{z-1/2} \sqrt{2\pi} \left(1 + \frac{1}{12z} + \frac{1}{288z^2} - \frac{139}{51840z^3} - \dots\right)$$

$$\Gamma\left(\frac{s}{2}\right) \pi^{-s/2} \zeta(s) = \Gamma\left(\frac{1-s}{2}\right) \pi^{-(1-s)/2} \zeta(1-s)$$

$$B(p, q) = \frac{\Gamma(p)\Gamma(q)}{\Gamma(p+q)} = \frac{(p-1)!(q-1)!}{(p+q-1)!}$$

$$\begin{aligned} m!n! &= \int_0^{\infty} e^{-u} u^m du \int_0^{\infty} e^{-v} v^n dv \quad \begin{matrix} u=x^2 \\ v=y^2 \end{matrix} \\ &= 4 \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-(x^2+y^2)} |x|^{2m+1} |y|^{2n+1} dx dy \quad \begin{matrix} x=r\cos\theta \\ y=r\sin\theta \end{matrix} \\ &= \int_0^{2\pi} \int_0^{\infty} e^{-r^2} |r\cos\theta|^{2m+1} |r\sin\theta|^{2n+1} r dr d\theta \\ &= \int_0^{\infty} e^{-r^2} r^{2m+2n+3} dr \int_0^{2\pi} |\cos\theta|^{2m+1} |\sin\theta|^{2n+1} d\theta \\ &= 4 \int_0^{\infty} e^{-r^2} r^{2m+2n+3} dr \int_0^{\pi/2} \cos^{2m+1}\theta \sin^{2n+1}\theta d\theta \\ &= 2(m+n+1)! \int_0^{\pi/2} \cos^{2m+1}\theta \sin^{2n+1}\theta d\theta \quad u=\cos^2\theta \end{aligned}$$

$$B(m+1, n+1) = \frac{m!n!}{(m+n+1)!} = \int_0^1 u^m (1-u)^n du$$

↳ can generalize to non-integers

↳ u-integrals easy to bound

$$\text{Use: } \left| \frac{\Gamma(a+iy)}{\Gamma(b+iy)} \right| = O_{a,b}(1) \quad \begin{matrix} \text{push in } \pi/2 \\ \text{push in } \pi/2 \end{matrix} \quad \frac{\Gamma(a+iy)\Gamma(b-a)}{\Gamma(b+iy)} = \int_0^1 t^{a+iy-1} (1+t)^{b-a-1} dt$$

↳ arises in shifts of L-fns

GAMMA AND BETA FUNCTIONS + DISTRIBUTIONS

GAMMA FUNCTION

$$\Gamma(\alpha) = \int_0^{\infty} y^{\alpha-1} e^{-y} dy \quad \text{for } \operatorname{Re}(\alpha) > 0$$

• NOTE: $\Gamma(1) = 1$ and $\Gamma(\alpha+1) = \alpha \Gamma(\alpha)$ for $\operatorname{Re}(\alpha) > 0$

$$\hookrightarrow \Gamma(1) = \int_0^{\infty} e^{-y} dy = 1$$

$$\Gamma(\alpha+1) = \int_0^{\infty} y^{\alpha} e^{-y} dy \quad \text{integrate by parts:}$$

$$\begin{aligned} & \begin{cases} u = y^{\alpha} & du = \alpha y^{\alpha-1} dy \\ dv = e^{-y} & v = -e^{-y} \end{cases} \\ & = -y^{\alpha} e^{-y} \Big|_0^{\infty} + \int_0^{\infty} \alpha y^{\alpha-1} e^{-y} dy \\ & = \alpha \int_0^{\infty} y^{\alpha-1} e^{-y} dy \\ & = \alpha \Gamma(\alpha) \end{aligned}$$

• NOTE: If $\alpha = n$ (a positive integer) then $\Gamma(n+1) = n!$
This follows from repeated use of $\Gamma(\alpha+1) = \alpha \Gamma(\alpha)$ and $\Gamma(1) = 1$

$$\text{THM: } \Gamma(1/2) = \sqrt{\pi}$$

we shall use this technique again in studying the Beta distribution.
Let $\alpha \in (0,1)$ so $\Gamma(\alpha), \Gamma(1-\alpha)$ exist and are finite. As all integrands below are positive, all change of variables can be justified.

$$\Gamma(\alpha) \Gamma(1-\alpha) = \int_0^{\infty} \int_0^{\infty} x^{\alpha-1} e^{-x} y^{1-\alpha-1} e^{-y} dx dy$$

$$= \int_0^{\infty} \int_0^{\infty} \left(\frac{x}{y}\right)^{\alpha-1} e^{-(\frac{x}{y})y} y^{-1} e^{-y} dx dy$$

$$\text{Change variables. Fix } y, \text{ set } z = \frac{x}{y}, dz = \frac{dx}{y} = y^{-1} dx$$

$$= \int_0^{\infty} \int_0^{\infty} z^{\alpha-1} e^{-zy} e^{-y} dz dy$$

$$= \int_0^{\infty} \int_0^{\infty} z^{\alpha-1} e^{-(z+1)y} dz dy \quad \begin{cases} \text{Change Vars: } u = (z+1)y, du = (z+1) dy \\ \text{Fix } z \end{cases}$$

$$= \int_0^{\infty} \int_0^{\infty} z^{\alpha-1} e^{-u} \frac{du}{z+1} dz = \int_0^{\infty} \frac{z^{\alpha-1}}{z+1} \underbrace{\int_0^{\infty} e^{-u} du}_{= \Gamma(1) = 1} dz$$

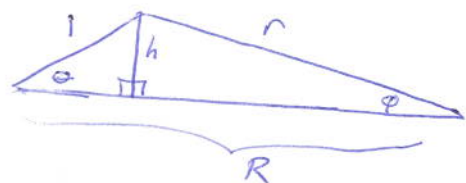
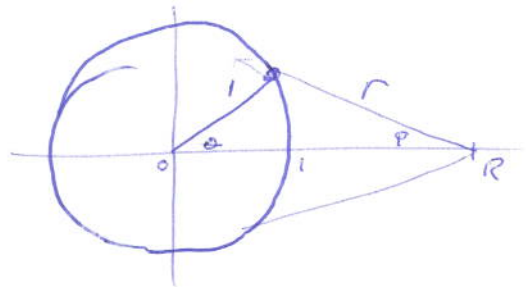
$$= \int_0^{\infty} \frac{z^{\alpha-1}}{1+z} dz \quad \text{Change vars: } z = u^2, dz = 2u du$$

$$= \int_0^{\infty} \frac{u^{2\alpha-2}}{(1+u^2)} \cdot 2u du \quad \text{Change vars: } u = \tan \theta, du = \sec^2 \theta d\theta$$

$$= 2 \int_0^{\pi/2} \frac{(\tan \theta)^{2\alpha-1}}{\sec^2 \theta} \sec^2 \theta d\theta \quad \begin{aligned} & 1 + \tan^2 \theta = \sec^2 \theta, u=0 \rightarrow \theta=0 \\ & u=\infty \rightarrow \theta=\pi/2 \end{aligned} \quad \text{get } 2 \int_0^{\pi/2} d\theta = \pi$$

6.3. Applications

Average values / center of mass: $\bar{x}_i = \frac{\int \dots \int x_i \delta(x_1, \dots, x_n) dx_1 \dots dx_n}{\int \dots \int \delta(x_1, \dots, x_n) dx_1 \dots dx_n}$



$$r^2 = R^2 + 1 - 2R \cos \theta$$

$$\text{Force from point: } \frac{G m_1 m_2}{r^2} \hat{e}_r$$

$$\text{Force } G \frac{m_1 m_2}{r^2} \cos \phi \cdot (-\hat{e}_x)$$

$$1 = R^2 + r^2 - 2Rr \cos \phi$$

$$\text{So } \cos \phi = \frac{R^2 + r^2 - 1}{2Rr}$$

$$\text{So } \cos \phi = \frac{2R^2 - 2R \cos \theta}{2Rr} = \frac{R - \cos \theta}{r}$$

$$\text{Force} = G m_1 m_2 \frac{R - \cos \theta}{(R^2 + 1 - 2R \cos \theta)^{3/2}}$$

$$\text{Force} = G m_1 m_2 \frac{(R - \cos \theta) \sin^3 \phi}{\sin^3 \theta}$$

$$\text{Force} = G m_1 m_2 \frac{R - \cos \theta}{(R^2 + 1 - 2R \cos \theta)^{3/2}} = G m_1 m_2 \frac{R - \cos \theta}{r^3}$$

$$\frac{r}{\sin \theta} = \frac{1}{\sin \phi}$$

$$r \sin \phi = \sin \theta$$

$$\frac{1}{r} = \frac{\sin \phi}{\sin \theta}$$

$$\sin \theta = r \sin \phi$$

$$r^2 = \frac{R^2 (R - \cos \theta)^2}{\cos^2 \phi}$$

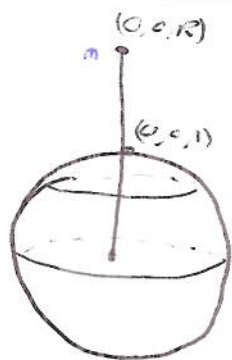
$$r = \frac{R(R - \cos \theta)}{\cos \phi}$$

Do in circular slices

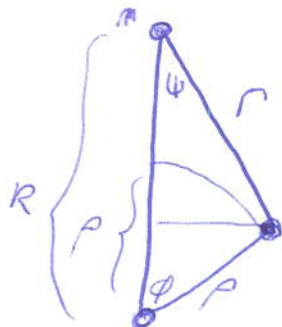
$$\hookrightarrow \text{Given } \theta, \text{ mass } 2\pi \sin \theta \cdot G m \frac{\cos \phi}{r^2} = 2\pi G m \frac{\sin \phi \cos \phi}{r}$$

Easier to do volumes! Cartesian lengths

6.3, APPLICATIONS



Assume unit density of 1, total mass $\frac{4}{3}\pi$



$$\cos \psi = \frac{R - \rho \cos \phi}{r} \quad r^2 = R^2 + \rho^2 - 2R\rho \cos \phi$$

Still true if $\phi > \pi/2$

F in $\hat{e}_r$ dir, amount in z -dir is $F \cos \psi$

$$\text{Force} = \iiint F(x,y,z) \cos \psi \, dx \, dy \, dz = \iiint F(\rho, \theta, \phi) \cos \psi \, \rho^2 \sin \phi \, d\rho \, d\theta \, d\phi$$

$$= Gm \iiint \frac{R - \rho \cos \phi}{(R^2 + \rho^2 - 2R\rho \cos \phi)^{3/2}} \rho^2 \sin \phi \, d\rho \, d\theta \, d\phi$$

$$= 2\pi Gm \int_{\rho=0}^1 \frac{\rho}{R} \int_{\phi=0}^{\pi} (R - \rho \cos \phi) \cdot (R^2 + \rho^2 - 2R\rho \cos \phi)^{-3/2} R \rho \sin \phi \, d\phi \, d\rho$$

$$u = R - \rho \cos \phi \quad du = (R^2 + \rho^2 - 2R\rho \cos \phi)^{-3/2} R \rho \sin \phi \, d\phi$$

$$du = \rho \sin \phi \, d\phi \quad v = -(R^2 + \rho^2 - 2R\rho \cos \phi)^{-1/2}$$

$$\text{Note } uv \Big|_0^{\pi} = \frac{R-\rho}{R-\rho} - \frac{R+\rho}{R+\rho} = 0$$

$$= 2\pi Gm \int_{\rho=0}^1 \frac{\rho}{R} \frac{1}{R} \int_{\phi=0}^{\pi} (R^2 + \rho^2 - 2R\rho \cos \phi)^{-1/2} \cdot R \rho \sin \phi \, d\phi$$

$$= \frac{2\pi Gm}{R^2} \int_{\rho=0}^1 \rho [(R+\rho) - (R-\rho)] \, d\rho$$

$$= Gm \left(\frac{4}{3}\pi \right) / R^2 \quad \text{as } \rho\text{-integral is } \int 2\rho^2 \, d\rho = \frac{2}{3}$$

See potential formulation in book is easier! Keep track of units

HW: Pg 404: #1

Suggested: # 3, #17, #18

Extra Credit: $F \propto 1/r^n$ in n -dim space: concentrated in center

6.4. IMPROPER IN. CORN

One variable: $\int_0^1 x^{-1/2} dx$ or $\int_1^\infty x^{-2} dx$

Use limits: $\lim_{\epsilon \rightarrow 0} \int_\epsilon^1 x^{-1/2} dx$

Several vars: look at contraction, expand

$$D = \{(x, y) : a \leq x \leq b, \phi_1(x) \leq y \leq \phi_2(x)\}$$

$$D_{\eta, \delta} = \{(x, y) : a + \eta \leq x \leq b - \eta, \phi_1(x) + \delta \leq y \leq \phi_2(x) - \delta\}$$

$$\lim_{(\eta, \delta) \rightarrow (0, 0)} \iint_{D_{\eta, \delta}} f dA$$

if this exists, then $\lim_{\eta \rightarrow 0} \lim_{\delta \rightarrow 0} \int_{a+\eta}^{b-\eta} \int_{\phi_1(x)+\delta}^{\phi_2(x)-\delta} f dA$ exists

However, iterated limits exist $\nrightarrow \lim_{(\eta, \delta) \rightarrow (0, 0)}$ exists

Example in book

Fubini: Definite, $f \geq 0$, $\left[\iint_D f dA, \int_a^b \int_{\phi_1(x)}^{\phi_2(x)} f dy dx \text{ for } y\text{-simple} \right.$
 $\left. \text{or } \int_a^b \int_{\phi_1(x)}^{\phi_2(x)} f dy dx \text{ for } x\text{-simple exist} \right]$ then f integrable
 and all equal.

Advanced: $\sum \sum a_{mn} \neq \sum \sum a_{mn}$

$$\begin{matrix} 0 & 0 \\ 0 & -1 \\ -1 & +1 \\ +1 & 0 \dots \end{matrix}$$

HW: pg 415: #5, #17, #18

Suggested: #1, #12, #13