

Math 383: Complex Analysis: Fall '21 (Williams)

Professor Steven J Miller: sjm1@williams.edu

Homepage:

[https://web.williams.edu/Mathematics/sjmiller/
public_html/383Fa21/](https://web.williams.edu/Mathematics/sjmiller/public_html/383Fa21/)

Lecture 30: 12-1-21: <https://youtu.be/mjl93gXaY8> ([slides](#))

Lecture 27: 11/15/17: Laplace's Method, Stirling's Formula: <https://youtu.be/AycMlf4Mbyo> (2015 lecture: https://youtu.be/GvKI5I_cfDQ)

Plan for the day: Lecture 30: December 1, 2021:

https://web.williams.edu/Mathematics/sjmiller/public_html/383Fa21/coursenotes/Math302_LecNotes_Intro.pdf

- Stationary Phase / Critical Points
- Dividing integer to maximize product
- Eigenvalues of Hermitian matrices

General items.

- Often easier to pass to a continuous analogue to study discrete problem
- Intuition from Taylor series: “higher” terms eventually negligible....

The Gamma function. The Gamma function $\Gamma(s)$ is

$$\Gamma(s) = \int_0^{\infty} e^{-x} x^{s-1} dx, \quad \Re(s) > 0.$$

$$x^{s-1} dx = x^s \frac{dx}{x}$$

Stirling's formula: As $n \rightarrow \infty$, we have

$$n! \approx n^n e^{-n} \sqrt{2\pi n};$$

by this we mean

$$\lim_{n \rightarrow \infty} \frac{n!}{n^n e^{-n} \sqrt{2\pi n}} = 1.$$

More precisely, we have the following series expansion:

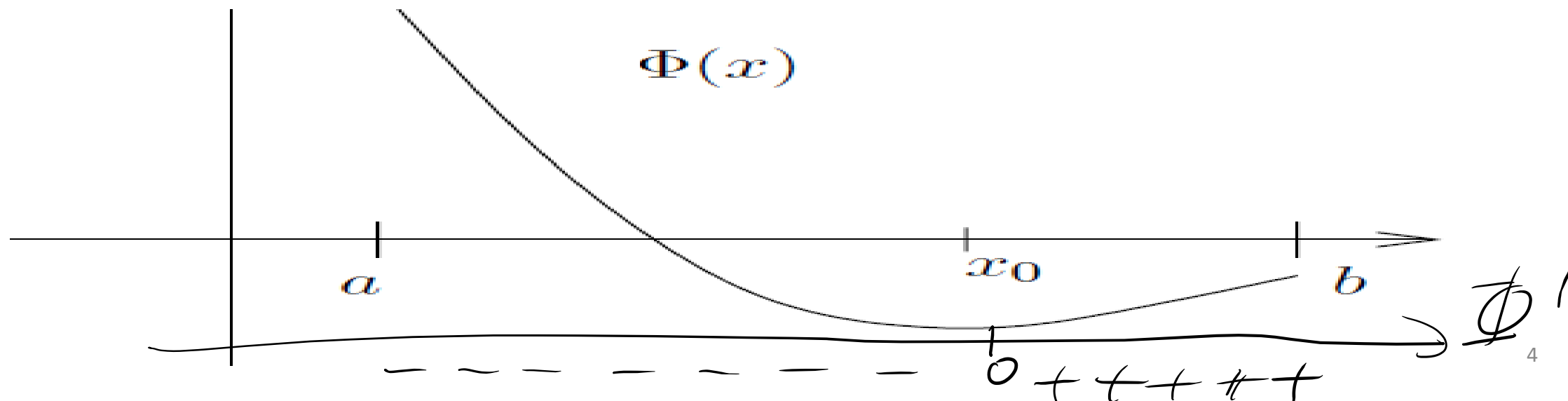
$$n! = n^n e^{-n} \sqrt{2\pi n} \left(1 + \frac{1}{12n} + \frac{1}{288n^2} - \frac{139}{51840n^3} - \cdots \right).$$

Consider

$$\underline{\Phi}(x) = \underline{\Phi}(x_0) + \underline{\Phi}'(x_0)(x-x_0) + \underline{\Phi}''(x_0)\frac{(x-x_0)^2}{2!} + O((x-x_0)^3)$$

$$\int_a^b e^{-s\Phi(x)} \psi(x) dx$$

where the **phase** Φ is real-valued, and both it and the **amplitude** ψ are assumed for simplicity to be indefinitely differentiable. Our hypothesis regarding the minimum of Φ is that there is an $x_0 \in (a, b)$ so that $\Phi'(x_0) = 0$, but $\Phi''(x) > 0$ throughout $[a, b]$ (Figure 2 illustrates the situation.)



Proposition 2.1 Under the above assumptions, with $s > 0$ and $s \rightarrow \infty$,

$$(8) \quad \int_a^b e^{-s\Phi(x)} \psi(x) dx = e^{-s\Phi(x_0)} \left[\frac{A}{s^{1/2}} + O\left(\frac{1}{s}\right) \right],$$

where

$$A = \sqrt{2\pi} \frac{\psi(x_0)}{(\Phi''(x_0))^{1/2}}.$$

Is this reasonable?

Φ is smallest at x_0 , have $e^{-s\Phi(x)}$
near x_0 , $\psi(x) \approx \psi(x_0)$

$\Phi(x)$ as $\Phi(x_0) + \Phi'(x_0)(x-x_0) + \frac{\Phi''(x_0)(x-x_0)^2}{2} + O((x-x_0)^3)$
after pull out $\Phi(x_0)$ have $\frac{1}{2}\Phi''(x_0)(x-x_0)^2 \rightarrow e^{-s\Phi''(x_0)(x-x_0)^2/2}$
looks like $N(x_0, 1/s\Phi''(x_0))$

Sketch:

$$\int_a^b \approx$$

$$\int_{x_0 - ?}^{x_0 + ?}$$

bounds depend on S

as $S \uparrow$, bounds $\downarrow$ (tighter about x_0)

Then

$$\int_{x_0 - \text{lower}}^{x_0 + \text{upper}} (\text{approx})$$

$$\approx \int_{-\infty}^{\infty} (\text{approx})$$

Wlog, $\Phi(x_0) = 0$, $\Phi(x) = \frac{\Phi''(x_0)(x-x_0)^2}{2!} + g(x)$

$$\int_a^b e^{-s\Phi(x)} \psi(x) dx = \int_{|x-x_0| < \delta(s)} + \int_{|x-x_0| > \delta(s)}$$

$\delta(s) = s^c$ or $\frac{1}{s^c}$
 c negative c positive,

or $\frac{1}{s^{1/2-c}} = \frac{s^c}{s^{1/2}}$

$a \leq x \leq b$
 Goes to zero as $s \rightarrow \infty$

$$e^{s\Phi(x_0)} \text{ vs } e^{s\Phi(x_0 + 1/\delta)} \approx e^{s\Phi(x_0) + \frac{1}{2}s\Phi''(x_0)/s}$$

$$e^{s\Phi(x_0)} \cdot e^{\Phi''(x_0)/2}$$

Take $\delta(s)$ to be $s^c/s^{1/2}$ with $0 < c < 1/2$

$$\int_a^{x_0 - \delta(s)} + \int_{x_0 + \delta(s)}^b e^{-s\Phi(x)} \psi(x) dx$$

$e^{-s\Phi(x)}$ largest at $x_0 - \delta(s)$ or $x_0 + \delta(s)$

$$e^{-s\left[\Phi(x_0) + \Phi''(x_0)\delta(s)^2/2 + O(\delta(s)^3)\right]}$$

$$= e^{-s\Phi(x_0)}$$

$$e^{-\frac{1}{2}\Phi''(x_0)s^{2c}}$$

Goes to zero
rapidly as
 $s \rightarrow \infty$

$$e^{O(s^{3c}/s^{1/2})}$$

need $3c < \frac{1}{2}$

or $c < 1/6$

$$\int_{x_0 - \delta(s)}^{x_0 + \delta(s)} e^{-s\Phi(x)} \psi(x) dx \quad \delta(s) = \frac{s^{-c}}{s^{1/2}} \quad 0 < c < \frac{1}{6}$$

as $\delta(s)$ so small as $s \rightarrow \infty$, $\psi(x) \approx \psi(x_0)$

$$e^{-s\Phi(x)} \approx \underbrace{e^{-s\Phi(x_0)}}_{\text{Fixed}} \underbrace{e^{-\frac{1}{2}s\Phi''(x_0)(x-x_0)^2}}_{\text{Gaussian}} \underbrace{e^{o((x-x_0)^3)}}_{(}$$

$$(\text{Integral}) \approx \int_{x_0 - \delta(s)}^{x_0 + \delta(s)} e^{-s\Phi(x_0)} e^{-\frac{(x-x_0)^2}{2} \frac{1}{s\Phi''(x_0)}} dx$$

as $s \rightarrow \infty$, get all the probability

Theorem 2.3 If $|s| \rightarrow \infty$ with $s \in S_\delta$, then

$$(11) \quad \Gamma(s) = e^{s \log s} e^{-s} \frac{\sqrt{2\pi}}{s^{1/2}} \left(1 + O \left(\frac{1}{|s|^{1/2}} \right) \right).$$

$$\Gamma(s) = \int_0^\infty e^{-x} x^s \frac{dx}{x} = \int_0^\infty e^{-x+s \log x} \frac{dx}{x}$$

rescale $x = sx$

$$\Gamma(s) = \int_0^\infty e^{-sx + s \log(sx)} \frac{1}{x} dx$$

using $\frac{dx}{x} \rightarrow \frac{dx}{x}$

$$= \int_0^\infty e^{-s(x - \log x)} e^{s \log s} \frac{1}{x} dx$$

note $e^{s \log s} = s^s$

$$= s^s \int_0^\infty e^{-s(x - \log x)} \frac{1}{x} dx$$

$$\Phi(x) = x - \log x$$

$$\Phi'(x) = 1 - 1/x, \quad \Phi'(1) = 0$$

$$\Phi''(x) = 1/x^2 > 0$$

$$\psi(x) = 1/x$$

Consider

$$\int_a^b e^{-s\Phi(x)} \psi(x) dx$$

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$$A = \sqrt{2\pi} \frac{\psi(x_0)}{(\Phi''(x_0))^{1/2}}.$$

LACOL DATA SCIENCE: Least Squares Lecture

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`http://www.williams.edu/Mathematics/sjmiller/
public_html/`

https://web.williams.edu/Mathematics/sjmiller/public_html/341Fa21/lectures/LACOL_DataScience_Regression1_Miller.pdf

<https://youtu.be/ZKoBBt933Ac>

Many non-linear relationships are linear after applying logarithms:

$$Y = BX^a \text{ then } \log(Y) = a \log(X) + b, \quad b = \log B.$$

Kepler's Third Law: if T is the orbital period of a planet traveling in an elliptical orbit about the sun (and no other objects exist), then $T^2 = \tilde{B}L^3$, where L is the length of the semi-major axis.

Assume do not know this – can we *discover* through statistics?

Logarithms and Applications

Many non-linear relationships are linear after applying logarithms:

$$Y = BX^a \text{ then } \log(Y) = a \log(X) + b, \quad b = \log B.$$

Kepler's Third Law: if T is the orbital period of a planet traveling in an elliptical orbit about the sun (and no other objects exist), then $T = BL^{1.5}$, where L is the length of the semi-major axis.

Assume do not know this – can we *discover* through statistics?

Kepler's Third Law: Can see the 1.5 exponent!

Data: Semi-major axis: Mercury 0.387, Venus 0.723, Earth 1.000, Mars 1.524, Jupiter 5.203, Saturn 9.539, Uranus 19.182, Neptune 30.06 (the units are astronomical units, where one astronomical unit is $1.496 \cdot 10^8$ km).

Data: orbital periods (in years) are 0.2408467, 0.61519726, 1.0000174, 1.8808476, 11.862615, 29.447498, 84.016846 and 164.79132.

If $T = BL^a$, what should B equal with this data? Units: bruno, millihelen, slug, smoot, See https://en.wikipedia.org/wiki/List_of_humorous_units_of_measurement

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Kepler's Third Law: Can see the 1.5 exponent!

If try $\log T = a \log L + b$: best fit values are...?

HOMEWORK!

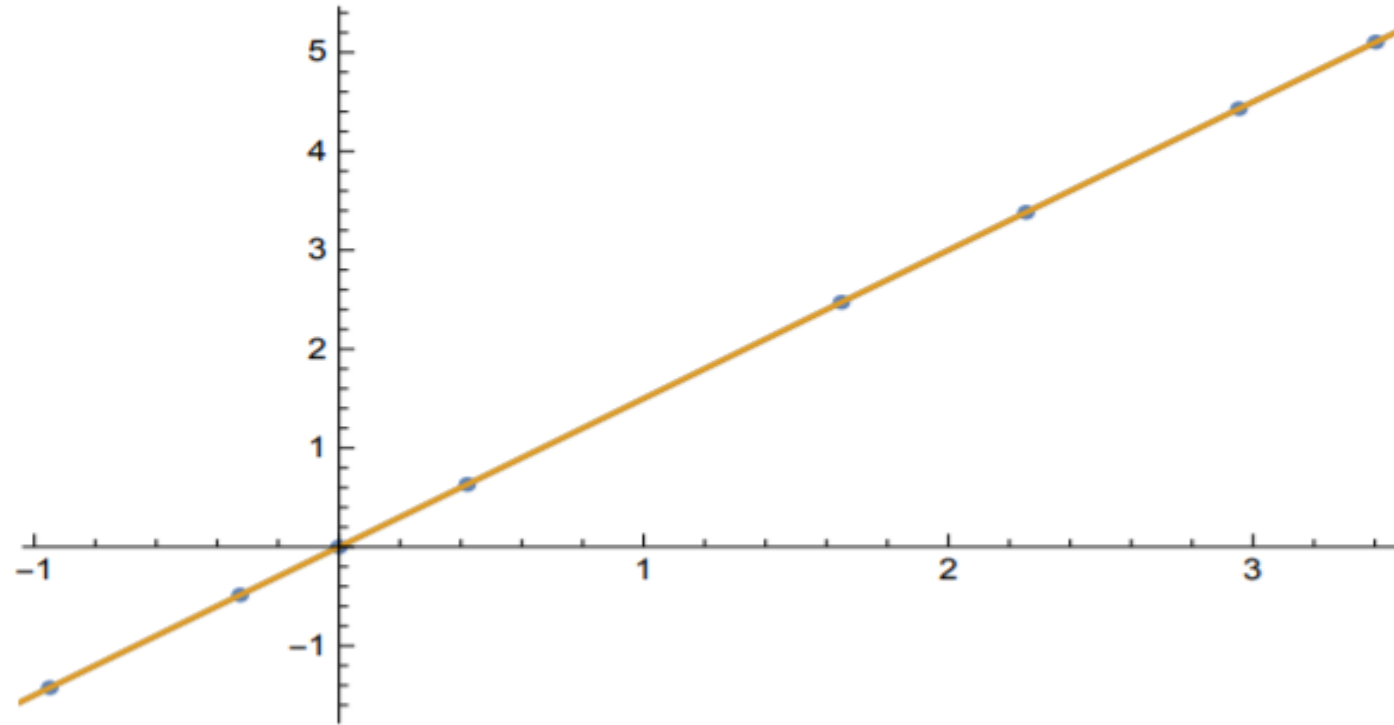


Figure: Plot of $\log P$ versus $\log L$ for planets. Is it surprising $b \approx 0$ (so $B \approx 1$ or $b \approx 0$)?

Word Counts

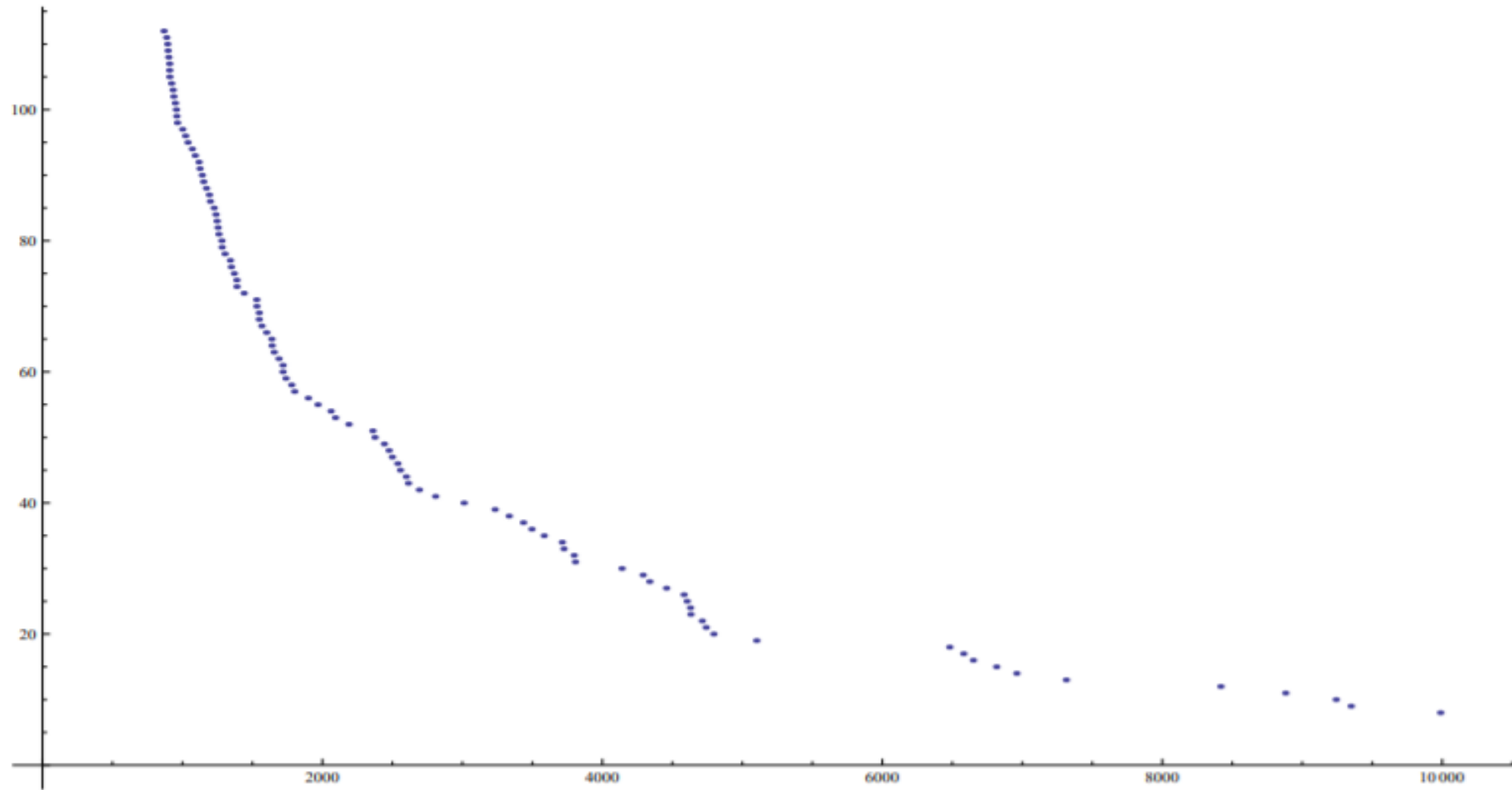


Figure: Plot of rank versus occurrences

Word Counts

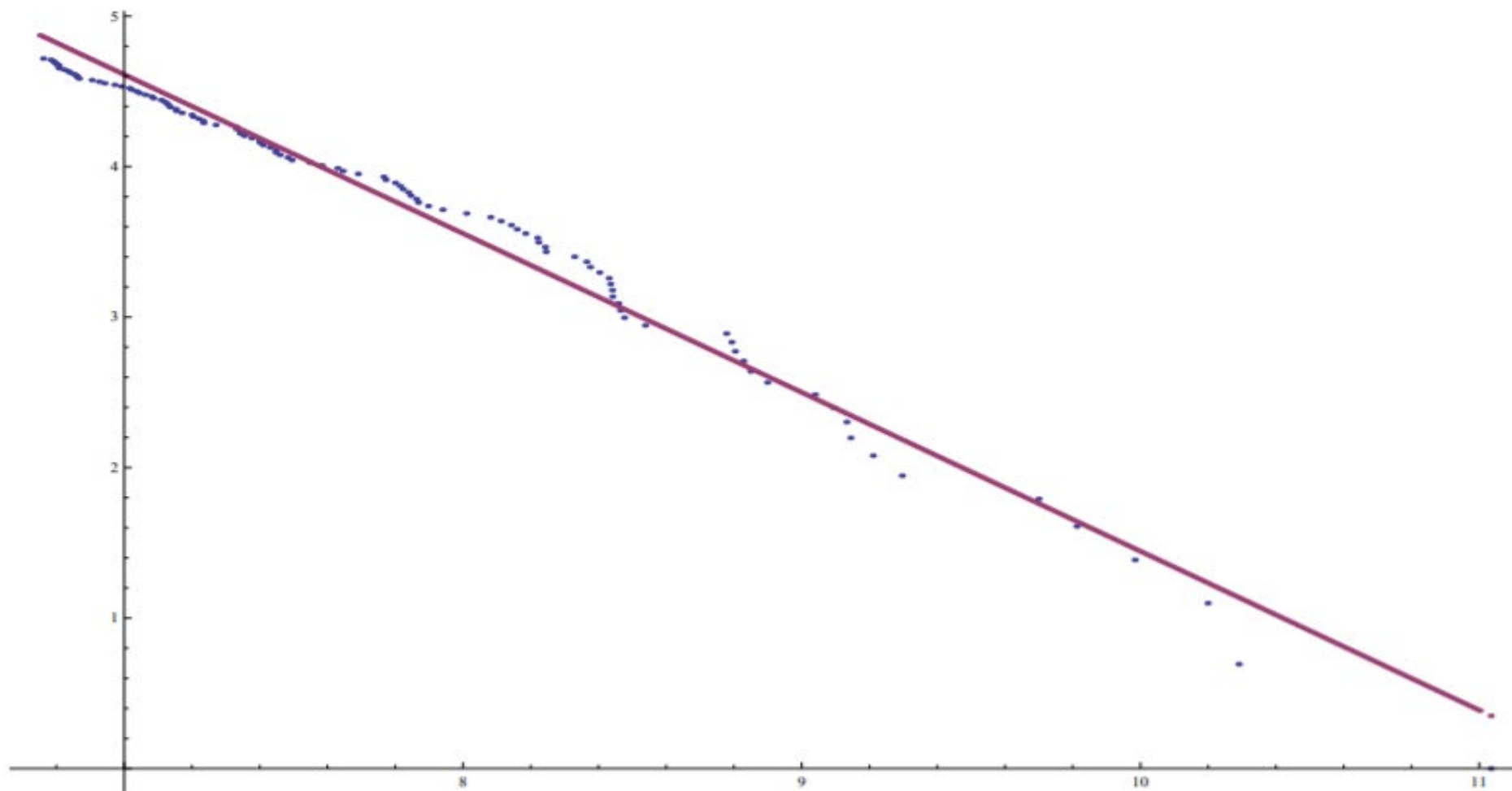
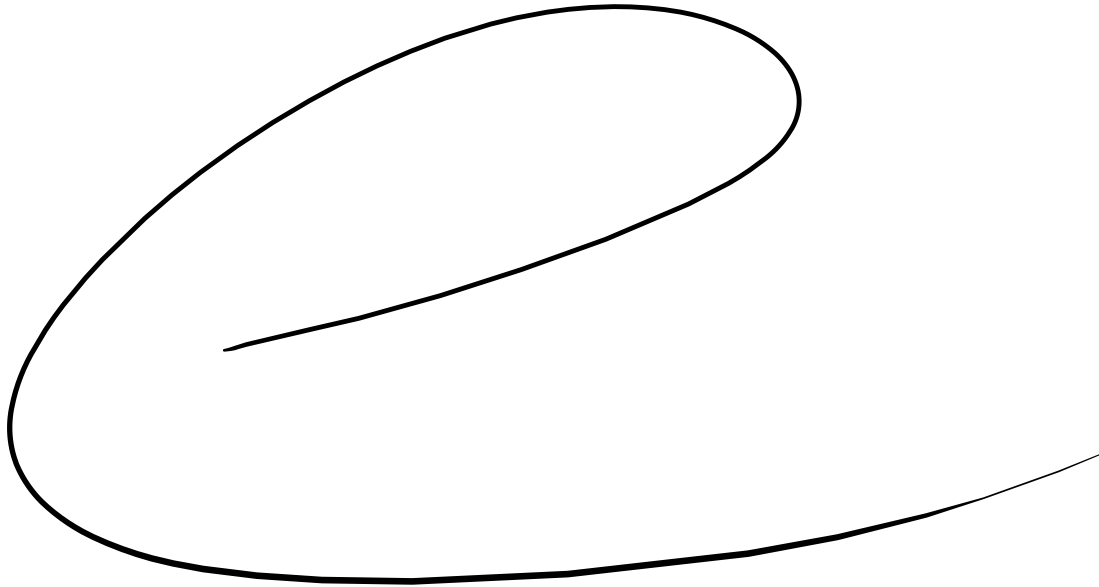


Figure: Plot of $\log(\text{rank})$ versus $\log(\text{occurrences})$

$$S = a_1 + \dots + a_n \quad \text{each } a_k \geq 1 \quad \max a_1, \dots, a_n$$

$$2 \leq a_k \leq 3$$

$$\underbrace{2+2+2}_8 \quad \text{vs} \quad \underbrace{3+3}_9$$



$$S = a_1 + \dots + a_n \quad \text{real} \quad 1 \leq a_k \leq 4$$

$$n=2: a_k = S/2$$

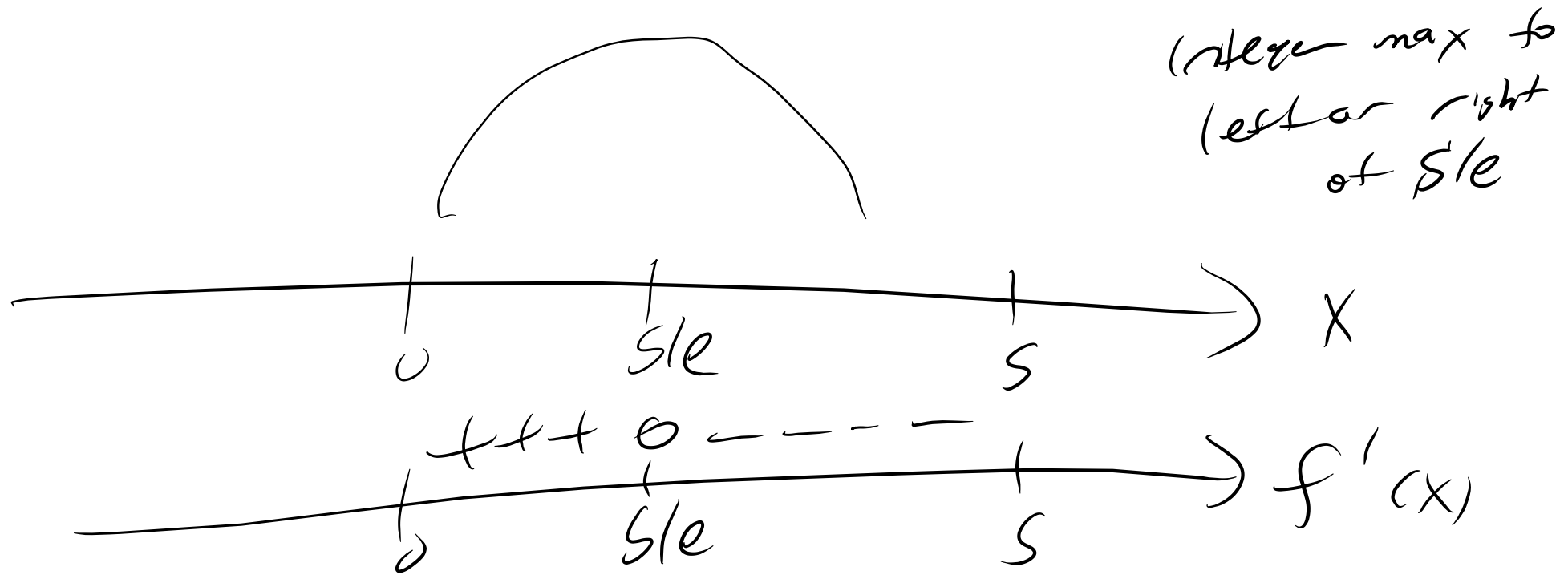
$$n=1 \quad a_k = S/1$$

$$f(n) = \left(\frac{S}{n}\right)^n \quad \text{max over } 1 \leq n \leq S$$

$$f(x) = \left(\frac{S}{x}\right)^x = e^{x \log(S/x)} = e^{x(\log S - \log x)}$$

$$f'(x) = f(x) \left[\log \frac{S}{x} + x \left(-\frac{1}{x}\right) \right] = f(x) \left[\log \left(\frac{S}{x}\right) - 1 \right]$$

$$= f(x) \log \left(\frac{S}{e^x}\right) = 0 \quad \text{if } \frac{S}{e^x} = 1 \Rightarrow x = S/e$$



$$f(x) = \left(\frac{S}{x}\right)^x$$

$$f'(x) = f(x) \left[\log \frac{S}{ex} \right]$$

CP was
 $x = S/e$

