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Random Matrix Ensembles with Split Limiting Behavior

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Introduction

Fundamental Problem: Spacing Between Events

General Formulation: Studying system, observe values at t_1, t_2, t_3, \ldots

Question: What rules govern the spacings between the t_i ?

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- Energy Levels of Nuclei.
- Eigenvalues of Matrices.
- Zeros of L-functions.
- Summands in Zeckendorf Decompositions.
- Primes.
- *n^k*α mod 1.

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Sketch of proofs

In studying many statistics, often three key steps:

- Determine correct scale for events.
- Oevelop an explicit formula relating what we want to study to something we understand.
- Use an averaging formula to analyze the quantities above.

It is not always trivial to figure out what is the correct statistic to study!



Background Material: Linear Algebra

Eigenvalue, Eigenvector

Say $\overrightarrow{v} \neq \overrightarrow{0}$ is an eigenvector of *A* with eigenvalue λ if $\overrightarrow{Av} = \lambda \overrightarrow{v}$.

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Say $\overrightarrow{v} \neq \overrightarrow{0}$ is an eigenvector of A with eigenvalue λ if $A\overrightarrow{v} = \lambda \overrightarrow{v}$.

Example:

$$\begin{pmatrix} 1 & 2 \\ 2 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = 3 \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$
$$\begin{pmatrix} 1 & 2 \\ 2 & 1 \end{pmatrix} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = -1 \begin{pmatrix} 1 \\ -1 \end{pmatrix}.$$

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Background Material: Probability

Probability Density

A random variable X has a probability density p(x) if

•
$$p(x) \ge 0;$$

•
$$\int_{-\infty}^{\infty} p(x) dx = 1;$$

•
$$\operatorname{Prob}(X \in [a, b]) = \int_a^b p(x) dx.$$

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Examples:

• Exponential:
$$p(x) = e^{-x/\lambda}/\lambda$$
 for $x \ge 0$;

2 Normal:
$$p(x) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-(x-\mu)^2/2\sigma^2};$$

Solution Uniform: $p(x) = \frac{1}{b-a}$ for $a \le x \le b$ and 0 otherwise.

Background Material: Probability (cont)

Key Concepts

• Mean (average value): $\mu = \int_{-\infty}^{\infty} xp(x) dx$.

Background Material: Probability (cont)

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- Variance (how spread out): $\sigma^2 = \int_{-\infty}^{\infty} (x \mu)^2 p(x) dx$.

Background Material: Probability (cont)

Key Concepts

- Mean (average value): $\mu = \int_{-\infty}^{\infty} xp(x) dx$.
- Variance (how spread out): $\sigma^2 = \int_{-\infty}^{\infty} (x \mu)^2 p(x) dx$.
- k^{th} moment: $\mu_k = \int_{-\infty}^{\infty} x^k p(x) dx$.

Background Material: Probability (cont)

Key Concepts

- Mean (average value): $\mu = \int_{-\infty}^{\infty} x p(x) dx$.
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•
$$k^{\text{th}}$$
 moment: $\mu_k = \int_{-\infty}^{\infty} x^k p(x) dx$.

Key observation

As a nice function is given by its Taylor series, a nice probability density is determined by its moments.

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Classical **Random Matrix Theory**

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Classical Mechanics: 3 Body Problem Intractable.

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Heavy nuclei (Uranium: 200+ protons / neutrons) worse!

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Get some info by shooting high-energy neutrons into nucleus, see what comes out.

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Fundamental Equation:

$$H\psi_n = E_n\psi_n$$

- H : matrix, entries depend on system
- E_n : energy levels
- ψ_n : energy eigenfunctions



Origins (continued)



- Statistical Mechanics: for each configuration, calculate quantity (say pressure).
- Average over all configurations most configurations close to system average.
- Nuclear physics: choose matrix at random, calculate eigenvalues, average over matrices (real Symmetric A = A^T, complex Hermitian A^T = A).

Random Matrix Ensembles

$$A = \begin{pmatrix} a_{11} & a_{12} & a_{13} & \cdots & a_{1N} \\ a_{12} & a_{22} & a_{23} & \cdots & a_{2N} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ a_{1N} & a_{2N} & a_{3N} & \cdots & a_{NN} \end{pmatrix} = A^{T}, \quad a_{ij} = a_{ji}$$

Random Matrix Ensembles

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Fix *p*, define

$$\mathsf{Prob}(A) = \prod_{1 \le i \le j \le N} p(a_{ij}).$$

Random Matrix Ensembles

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Fix *p*, define

$$\mathsf{Prob}(A) = \prod_{1 \le i \le j \le N} p(a_{ij}).$$

This means

$$\operatorname{Prob}\left(\boldsymbol{A}: \boldsymbol{a}_{ij} \in [\alpha_{ij}, \beta_{ij}]\right) = \prod_{1 \leq i \leq j \leq N} \int_{\boldsymbol{x}_{ij}=\alpha_{ij}}^{\beta_{ij}} \boldsymbol{p}(\boldsymbol{x}_{ij}) d\boldsymbol{x}_{ij}.$$

Eigenvalue Trace Lemma

Want to understand the eigenvalues of *A*, but it is the matrix elements that are chosen randomly and independently.

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Eigenvalue Trace Lemma

Let *A* be an $N \times N$ matrix with eigenvalues $\lambda_i(A)$. Then

Trace
$$(\mathbf{A}^k) = \sum_{n=1}^N \lambda_i(\mathbf{A})^k$$
,

where

Trace(
$$A^k$$
) = $\sum_{i_1=1}^{N} \cdots \sum_{i_k=1}^{N} a_{i_1 i_2} a_{i_2 i_3} \cdots a_{i_N i_1}$.



$$\delta(x - x_0)$$
 is a unit point mass at x_0 :
 $\int_{-\infty}^{\infty} f(x)\delta(x - x_0)dx = f(x_0).$



Eigenvalue Distribution

$$\delta(x - x_0)$$
 is a unit point mass at x_0 :
 $\int_{-\infty}^{\infty} f(x)\delta(x - x_0)dx = f(x_0).$

To each A, attach a probability measure:

$$\mu_{A,N}(x) = \frac{1}{N} \sum_{i=1}^{N} \delta\left(x - \frac{\lambda_i(A)}{2\sqrt{N}}\right)$$
$$\int_a^b \mu_{A,N}(x) dx = \frac{\#\left\{\lambda_i : \frac{\lambda_i(A)}{2\sqrt{N}} \in [a, b]\right\}}{N}$$
$$k^{\text{th}} \text{ moment} = \frac{\sum_{i=1}^{N} \lambda_i(A)^k}{2^k N^{\frac{k}{2}+1}} = \frac{\text{Trace}(A^k)}{2^k N^{\frac{k}{2}+1}}.$$

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Density of States

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Wigner's Semi-Circle Law

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Wigner's Semi-Circle Law

 $N \times N$ real symmetric matrices, entries i.i.d.r.v. from a fixed p(x) with mean 0, variance 1, and other moments finite. Then for almost all *A*, as $N \rightarrow \infty$

$$\mu_{A,N}(x) \longrightarrow egin{cases} rac{2}{\pi}\sqrt{1-x^2} & ext{if } |x| \leq 1 \ 0 & ext{otherwise.} \end{cases}$$

SKETCH OF PROOF: Correct Scale

Trace(
$$A^2$$
) = $\sum_{i=1}^N \lambda_i(A)^2$.

By the Central Limit Theorem:

$$Trace(A^{2}) = \sum_{i=1}^{N} \sum_{j=1}^{N} a_{ij}a_{ji} = \sum_{i=1}^{N} \sum_{j=1}^{N} a_{ij}^{2} \sim N^{2}$$
$$\sum_{i=1}^{N} \lambda_{i}(A)^{2} \sim N^{2}$$

Gives NAve $(\lambda_i(A)^2) \sim N^2$ or Ave $(\lambda_i(A)) \sim \sqrt{N}$.

SKETCH OF PROOF: Averaging Formula

Recall *k*-th moment of $\mu_{A,N}(x)$ is $\operatorname{Trace}(A^k)/2^k N^{k/2+1}$.

Average *k*-th moment is

$$\int \cdots \int \frac{\operatorname{Trace}(A^k)}{2^k N^{k/2+1}} \prod_{i \leq j} p(a_{ij}) da_{ij}.$$

Proof by method of moments: Two steps

- Show average of *k*-th moments converge to moments of semi-circle as *N* → ∞;
- Control variance (show it tends to zero as $N \to \infty$).

SKETCH OF PROOF: Averaging Formula for Second Moment

Substituting into expansion gives

$$\frac{1}{2^2N^2}\int_{-\infty}^{\infty}\cdots\int_{-\infty}^{\infty}\sum_{i=1}^{N}\sum_{j=1}^{N}a_{ij}^2\cdot p(a_{11})da_{11}\cdots p(a_{NN})da_{NN}$$

Integration factors as

$$\int_{a_{ij}=-\infty}^{\infty}a_{ij}^2p(a_{ij})da_{ij} \cdot \prod_{(k,l)\neq(i,j)\atop k< l}\int_{a_{kl}=-\infty}^{\infty}p(a_{kl})da_{kl} = 1.$$

Higher moments involve more advanced combinatorics (Catalan numbers).

Numerical example: Gaussian density



Numerical example: Cauchy density $p(x) = 1/(\pi(1 + x^2))$

Numerical example: Cauchy density $p(x) = 1/(\pi(1 + x^2))$



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Real Symmetric Toeplitz Matrices Chris Hammond and Steven J. Miller
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Toeplitz Ensembles

Toeplitz matrix is of the form

$$\begin{pmatrix} b_0 & b_1 & b_2 & \cdots & b_{N-1} \\ b_{-1} & b_0 & b_1 & \cdots & b_{N-2} \\ b_{-2} & b_{-1} & b_0 & \cdots & b_{N-3} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ b_{1-N} & b_{2-N} & b_{3-N} & \cdots & b_0 \end{pmatrix}$$

- Will consider Real Symmetric Toeplitz matrices.
- Main diagonal zero, N 1 independent parameters.
- Normalize Eigenvalues by \sqrt{N} .

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Numerical Observations: Thoughts?



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Eigenvalue Density Measure

$$\mu_{A,N}(x)dx = \frac{1}{N}\sum_{i=1}^{N}\delta\left(x-\frac{\lambda_{i}(A)}{\sqrt{N}}\right)dx.$$

The k^{th} moment of $\mu_{A,N}(x)$ is

$$M_k(A, N) = \frac{1}{N^{\frac{k}{2}+1}} \sum_{i=1}^N \lambda_i^k(A) = \frac{\operatorname{Trace}(A^k)}{N^{\frac{k}{2}+1}}.$$

Let

$$M_k = \lim_{N \to \infty} \mathbb{E}_A [M_k(A, N)];$$

have $M_2 = 1$ and $M_{2k+1} = 0$.

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Even Moments

$$M_{2k}(N) = \frac{1}{N^{k+1}} \sum_{1 \leq i_1, \cdots, i_{2k} \leq N} \mathbb{E}(b_{|i_1 - i_2|} b_{|i_2 - i_3|} \cdots b_{|i_{2k} - i_1|}).$$

Main Term: b_i 's matched in pairs, say

 $b_{|i_m-i_{m+1}|} = b_{|i_n-i_{n+1}|}, \quad x_m = |i_m-i_{m+1}| = |i_n-i_{n+1}|.$

Two possibilities:

$$i_m - i_{m+1} = i_n - i_{n+1}$$
 or $i_m - i_{m+1} = -(i_n - i_{n+1})$.

(2k-1)!! ways to pair, 2^k choices of sign.

Main Term: All Signs Negative (else lower order contribution)

$$M_{2k}(N) = \frac{1}{N^{k+1}} \sum_{1 \leq i_1, \cdots, i_{2k} \leq N} \mathbb{E}(b_{|i_1 - i_2|} b_{|i_2 - i_3|} \cdots b_{|i_{2k} - i_1|}).$$

Let x_1, \ldots, x_k be the values of the $|i_j - i_{j+1}|$'s, $\epsilon_1, \ldots, \epsilon_k$ the choices of sign. Define $\widetilde{x}_1 = i_1 - i_2$, $\widetilde{x}_2 = i_2 - i_3$,

$$i_2 = i_1 - \widetilde{x}_1$$

$$i_3 = i_1 - \widetilde{x}_1 - \widetilde{x}_2$$

$$\vdots$$

$$i_1 = i_1 - \widetilde{x}_1 - \cdots - \widetilde{x}_{2k}$$

$$\widetilde{x}_1 + \cdots + \widetilde{x}_{2k} = \sum_{j=1}^k (1 + \epsilon_j) \eta_j x_j = 0, \quad \eta_j = \pm 1.$$

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Even Moments: Summary

Main Term: paired, all signs negative.

$$M_{2k}(N) \leq (2k-1)!! + O_k\left(\frac{1}{N}x\right).$$

Bounded by Gaussian.

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The Fourth Moment



$$M_4(N) = \frac{1}{N^3} \sum_{1 \le i_1, i_2, i_3, i_4 \le N} \mathbb{E}(b_{|i_1 - i_2|} b_{|i_2 - i_3|} b_{|i_3 - i_4|} b_{|i_4 - i_1|})$$

Let
$$x_j = |i_j - i_{j+1}|$$
.

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The Fourth Moment

Case One: $x_1 = x_2, x_3 = x_4$:

$$i_1 - i_2 = -(i_2 - i_3)$$
 and $i_3 - i_4 = -(i_4 - i_1)$.

Implies

 $i_1 = i_3$, i_2 and i_4 arbitrary.

Left with $\mathbb{E}[b_{x_1}^2 b_{x_3}^2]$: $N^3 - N$ times get 1, N times get $p_4 = \mathbb{E}[b_{x_1}^4]$. Contributes 1 in the limit. Introduction Classical RMT Density of States Concerned to the concerned to

The Fourth Moment

$$M_4(N) = \frac{1}{N^3} \sum_{1 \le i_1, i_2, i_3, i_4 \le N} \mathbb{E}(b_{|i_1 - i_2|} b_{|i_2 - i_3|} b_{|i_3 - i_4|} b_{|i_4 - i_1|})$$

Case Two: Diophantine Obstruction: $x_1 = x_3$ and $x_2 = x_4$.

$$i_1 - i_2 = -(i_3 - i_4)$$
 and $i_2 - i_3 = -(i_4 - i_1)$.

This yields

$$i_1 = i_2 + i_4 - i_3, \ i_1, i_2, i_3, i_4 \in \{1, \dots, N\}.$$

If i_2 , $i_4 \ge \frac{2N}{3}$ and $i_3 < \frac{N}{3}$, $i_1 > N$: at most $(1 - \frac{1}{27})N^3$ valid choices.



The Fourth Moment

Theorem: Fourth Moment: Let p_4 be the fourth moment of *p*. Then

$$M_4(N) = 2\frac{2}{3} + O_{p_4}\left(\frac{1}{N}\right).$$

500 Toeplitz Matrices, 400×400 .



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Main Result

Theorem: HM '05

For real symmetric Toeplitz matrices, the limiting spectral measure converges in probability to a unique measure of unbounded support which is not the Gaussian. If p is even have strong convergence).

Massey, Miller and Sinsheimer '07 proved that if first row is a palindrome converges to a Gaussian.

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Block Circulant Ensemble

With Murat Koloğlu, Gene Kopp, Fred Strauch and Wentao Xiong.

The Ensemble of *m*-Block Circulant Matrices

Symmetric matrices periodic with period *m* on wrapped diagonals, i.e., symmetric block circulant matrices.

8-by-8 real symmetric 2-block circulant matrix:

$$\begin{pmatrix} c_0 & c_1 & c_2 & c_3 & c_4 & d_3 & c_2 & d_1 \\ c_1 & d_0 & d_1 & d_2 & d_3 & d_4 & c_3 & d_2 \\ \hline c_2 & d_1 & c_0 & c_1 & c_2 & c_3 & c_4 & d_3 \\ \hline c_3 & d_2 & c_1 & d_0 & d_1 & d_2 & d_3 & d_4 \\ \hline c_4 & d_3 & c_2 & d_1 & c_0 & c_1 & c_2 & c_3 \\ \hline d_3 & d_4 & c_3 & d_2 & c_1 & d_0 & d_1 & d_2 \\ \hline c_2 & c_3 & c_4 & d_3 & c_2 & d_1 & c_0 & c_1 \\ \hline d_1 & d_2 & d_3 & d_4 & c_3 & d_2 & c_1 & d_0 \end{pmatrix}$$

Choose distinct entries i.i.d.r.v.

Oriented Matchings and Dualization

Compute moments of eigenvalue distribution (as *m* stays fixed and $N \rightarrow \infty$) using the combinatorics of pairings. Rewrite:

$$M_n(N) = \frac{1}{N^{\frac{n}{2}+1}} \sum_{1 \le i_1, \dots, i_n \le N} \mathbb{E}(a_{i_1 i_2} a_{i_2 i_3} \cdots a_{i_n i_1})$$

= $\frac{1}{N^{\frac{n}{2}+1}} \sum_{\sim} \eta(\sim) m_{d_1(\sim)} \cdots m_{d_l(\sim)}.$

where the sum is over oriented matchings on the edges $\{(1,2), (2,3), ..., (n,1)\}$ of a regular *n*-gon.

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Oriented Matchings and Dualization



Figure: An oriented matching in the expansion for $M_n(N) = M_6(8)$.



Contributing Terms

As $N \to \infty$, the only terms that contribute to this sum are those in which the entries are matched in pairs and with opposite orientation.





Only Topology Matters

Think of pairings as topological identifications; the contributing ones give rise to orientable surfaces.



Contribution from such a pairing is m^{-2g} , where *g* is the genus (number of holes) of the surface. Proof: combinatorial argument involving Euler characteristic.

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Computing the Even Moments

Theorem: Even Moment Formula

$$M_{2k} = \sum_{g=0}^{\lfloor k/2
floor} \varepsilon_g(k) m^{-2g} + O_k\left(rac{1}{N}
ight),$$

with $\varepsilon_g(k)$ the number of pairings of the edges of a (2k)-gon giving rise to a genus *g* surface.

J. Harer and D. Zagier (1986) gave generating functions for the $\varepsilon_g(k)$.

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Harer and Zagier

$$\sum_{g=0}^{\lfloor k/2 \rfloor} \varepsilon_g(k) r^{k+1-2g} = (2k-1)!! c(k,r)$$

where

$$1 + 2\sum_{k=0}^{\infty} c(k,r) x^{k+1} = \left(\frac{1+x}{1-x}\right)^{r}$$

Thus, we write

$$M_{2k} = m^{-(k+1)}(2k-1)!! c(k,m).$$

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A multiplicative convolution and Cauchy's residue formula yield the characteristic function of the distribution.

$$\begin{split} \phi(t) &= \sum_{k=0}^{\infty} \frac{(it)^{2k} M_{2k}}{(2k)!} = \frac{1}{m} \sum_{k=0}^{\infty} \frac{(-t^2/2m)^k}{k!} c(k,m) \\ &= \frac{1}{2\pi i m} \oint_{|z|=2} \frac{1}{2z^{-1}} \left(\left(\frac{1+z^{-1}}{1-z^{-1}} \right)^m - 1 \right) e^{-t^2 z/2m} \frac{dz}{z} \\ &= \frac{1}{m} e^{\frac{-t^2}{2m}} \sum_{\ell=1}^m \binom{m}{\ell} \frac{1}{(\ell-1)!} \left(\frac{-t^2}{m} \right)^{\ell-1}. \end{split}$$

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Results

Fourier transform and algebra yields

Theorem: Koloğlu, Kopp and Miller

The limiting spectral density function $f_m(x)$ of the real symmetric *m*-block circulant ensemble is given by the formula

$$f_m(x) = \frac{e^{-\frac{mx^2}{2}}}{\sqrt{2\pi m}} \sum_{r=0}^m \frac{1}{(2r)!} \sum_{s=0}^{m-r} \binom{m}{r+s+1} \frac{(2r+2s)!}{(r+s)!s!} \left(-\frac{1}{2}\right)^s (mx^2)^r.$$

As $m \to \infty$, the limiting spectral densities approach the semicircle distribution.

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Results (continued)



Figure: Plot for f_1 and histogram of eigenvalues of 100 circulant matrices of size 400×400 .





Figure: Plot for f_2 and histogram of eigenvalues of 100 2-block circulant matrices of size 400 × 400.





Figure: Plot for f_3 and histogram of eigenvalues of 100 3-block circulant matrices of size 402 × 402.





Figure: Plot for f_4 and histogram of eigenvalues of 100 4-block circulant matrices of size 400 × 400.





Figure: Plot for f_8 and histogram of eigenvalues of 100 8-block circulant matrices of size 400 × 400.

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Results (continued)



Figure: Plot for f_{20} and histogram of eigenvalues of 100 20-block circulant matrices of size 400 × 400.





Figure: Plot of convergence to the semi-circle.

The Limiting Spectral Measure for Ensembles of Symmetric Block Circulant Matrices (with Murat Koloğlu, Gene S. Kopp, Frederick W. Strauch and Wentao Xiong), Journal of Theoretical Probability **26** (2013), no. 4, 1020–1060. http://arxiv.org/abs/1008.4812 Introduction Classical RMT Density of States cooperative States cooper

Checkerboard Matrices

- First paper with Paula Burkhardt, Peter Cohen, Jonathan Dewitt, Max Hlavacek, Carsten Sprunger, Yen Nhi Truong Vu, Roger Van Peski, and Kevin Yang, and an appendix joint with Manuel Fernandez and Nicholas Sieger.
- Second paper with Ryan Chen, Yujin Kim, Jared Lichtman, Shannon Sweitzer, and Eric Winsor.
- Third paper with Fangu Chen, Yuxin Lin and Jiahui Yu.

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Checkerboard Matrices: $N \times N(k, w)$ -checkerboard ensemble

Matrices $M = (m_{ij}) = M^T$ with a_{ij} iidrv, mean 0, variance 1, finite higher moments, *w* fixed and

$$m_{ij} = egin{cases} {a_{ij}} & ext{if } i
eq j \mod k \ w & ext{if } i \equiv j \mod k. \end{cases}$$

Example: (3, w)-checkerboard matrix:

$$\begin{pmatrix} \mathbf{W} & \mathbf{a}_{0,1} & \mathbf{a}_{0,2} & \mathbf{W} & \mathbf{a}_{0,4} & \cdots & \mathbf{a}_{0,N-1} \\ \mathbf{a}_{1,0} & \mathbf{W} & \mathbf{a}_{1,2} & \mathbf{a}_{1,3} & \mathbf{W} & \cdots & \mathbf{a}_{1,N-1} \\ \mathbf{a}_{2,0} & \mathbf{a}_{2,1} & \mathbf{W} & \mathbf{a}_{2,3} & \mathbf{a}_{2,4} & \cdots & \mathbf{W} \\ \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ \mathbf{a}_{0,N-1} & \mathbf{a}_{1,N-1} & \mathbf{W} & \mathbf{a}_{3,N-1} & \mathbf{a}_{4,N-1} & \cdots & \mathbf{W} \end{pmatrix}$$

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Split Eigenvalue Distribution



Figure: Histogram of normalized eigenvalues: 2-checkerboard 100×100 matrices, 100 trials.

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Split Eigenvalue Distribution



Figure: Histogram of normalized eigenvalues: 2-checkerboard 150 \times 150 matrices, 100 trials.

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Split Eigenvalue Distribution



Figure: Histogram of normalized eigenvalues: 2-checkerboard 200×200 matrices, 100 trials.

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Split Eigenvalue Distribution



Figure: Histogram of normalized eigenvalues: 2-checkerboard 250×250 matrices, 100 trials.

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Split Eigenvalue Distribution



Figure: Histogram of normalized eigenvalues: 2-checkerboard 300×300 matrices, 100 trials.



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Split Eigenvalue Distribution



Figure: Histogram of normalized eigenvalues: 2-checkerboard 350×350 matrices, 100 trials.


The Weighting Function

Use weighting function $f_n(x) = x^{2n}(x-2)^{2n}$.



Figure: $f_n(x)$ plotted for $n \in \{1, 2, 3, 4\}$.



The Weighting Function

Use weighting function $f_n(x) = x^{2n}(x-2)^{2n}$.



Figure: $f_n(x)$ plotted for $n = 4^m, m \in \{0, 1, ..., 5\}$.

Spectral distribution of hollow GOE



Figure: Hist. of eigenvals of 32000 (Left) 2×2 hollow GOE matrices, (Right) 3×3 hollow GOE matrices.



Figure: Hist. of eigenvals of 32000 (Left) 4×4 hollow GOE matrices, (Right) 16×16 hollow GOE matrices.

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Riemann Zeta Function

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s} = \prod_{p \text{ prime}} \left(1 - \frac{1}{p^s}\right)^{-1}, \quad \text{Re}(s) > 1.$$

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Unique Factorization: $n = p_1^{r_1} \cdots p_m^{r_m}$.

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Unique Factorization: $n = p_1^{r_1} \cdots p_m^{r_m}$.

$$\prod_{p} \left(1 - \frac{1}{p^{s}}\right)^{-1} = \left[1 + \frac{1}{2^{s}} + \left(\frac{1}{2^{s}}\right)^{2} + \cdots\right] \left[1 + \frac{1}{3^{s}} + \left(\frac{1}{3^{s}}\right)^{2} + \cdots\right] \cdots$$
$$= \sum_{n} \frac{1}{n^{s}}.$$

Riemann Zeta Function (cont)

$$\begin{aligned} \zeta(s) &= \sum_{n} \frac{1}{n^{s}} = \prod_{p} \left(1 - \frac{1}{p^{s}} \right)^{-1}, \quad \operatorname{Re}(s) > 1 \\ \pi(x) &= \#\{p : p \text{ is prime}, p \le x\} \end{aligned}$$

Properties of $\zeta(s)$ and Primes:

Riemann Zeta Function (cont)

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Properties of $\zeta(s)$ and Primes:

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$$\lim_{s \to 1^+} \zeta(s) = \infty, \pi(x) \to \infty.$$

Riemann Zeta Function (cont)

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Properties of $\zeta(s)$ and Primes:

•
$$\lim_{s \to 1^+} \zeta(s) = \infty, \pi(x) \to \infty.$$

• $\zeta(2) = \frac{\pi^2}{6}, \pi(x) \to \infty.$

Riemann Zeta Function

$$\zeta(s) = \sum_{n=1}^{\infty} \frac{1}{n^s} = \prod_{p \text{ prime}} \left(1 - \frac{1}{p^s}\right)^{-1}, \quad \text{Re}(s) > 1.$$

Functional Equation:

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$$\xi(s) = \Gamma\left(\frac{s}{2}\right)\pi^{-\frac{s}{2}}\zeta(s) = \xi(1-s).$$

Riemann Hypothesis (RH):

All non-trivial zeros have $\operatorname{Re}(s) = \frac{1}{2}$; can write zeros as $\frac{1}{2} + i\gamma$.

Observation: Spacings b/w zeros appear same as b/w eigenvalues of Complex Hermitian matrices $\overline{A}^T = A$.



General *L*-functions

$$L(s, f) = \sum_{n=1}^{\infty} \frac{a_f(n)}{n^s} = \prod_{p \text{ prime}} L_p(s, f)^{-1}, \quad \text{Re}(s) > 1.$$

Functional Equation:

$$\Lambda(\boldsymbol{s},f) = \Lambda_{\infty}(\boldsymbol{s},f)L(\boldsymbol{s},f) = \Lambda(1-\boldsymbol{s},f).$$

Generalized Riemann Hypothesis (RH):

All non-trivial zeros have $\operatorname{Re}(s) = \frac{1}{2}$; can write zeros as $\frac{1}{2} + i\gamma$.

Observation: Spacings b/w zeros appear same as b/w eigenvalues of Complex Hermitian matrices $\overline{A}^T = A$.

Nuclear spacings: Thorium



227 spacings b/w adjacent energy levels of Thorium.



Zeros of $\zeta(s)$ vs GUE



70 million spacings b/w adjacent zeros of $\zeta(s)$, starting at the 10^{20th} zero (from Odlyzko).

Elliptic Curves: Mordell-Weil Group

Elliptic curve $y^2 = x^3 + ax + b$ with rational solutions $P = (x_1, y_1)$ and $Q = (x_2, y_2)$ and connecting line y = mx + b.





Adding a point P to itself

Addition of distinct points P and Q

 $E(\mathbb{Q}) \approx E(\mathbb{Q})_{\mathrm{tors}} \oplus \mathbb{Z}^{r}$



Elliptic curve *L*-function

 $E: y^2 = x^3 + ax + b$, associate *L*-function

$$L(s,E) = \sum_{n=1}^{\infty} \frac{a_E(n)}{n^s} = \prod_{p \text{ prime}} L_E(p^{-s}),$$

where

$$a_{\mathcal{E}}(p) = p - \#\{(x, y) \in (\mathbb{Z}/p\mathbb{Z})^2 : y^2 \equiv x^3 + ax + b \bmod p\}.$$



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Birch and Swinnerton-Dyer Conjecture

Rank of group of rational solutions equals order of vanishing of L(s, E) at s = 1/2.

Properties of zeros of *L*-functions

- infinitude of primes, primes in arithmetic progression.
- Chebyshev's bias: $\pi_{3,4}(x) \ge \pi_{1,4}(x)$ 'most' of the time.
- Birch and Swinnerton-Dyer conjecture.
- Goldfeld, Gross-Zagier: bound for *h*(*D*) from *L*-functions with many central point zeros.
- Even better estimates for *h*(*D*) if a positive percentage of zeros of ζ(*s*) are at most 1/2 − ε of the average spacing to the next zero.

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Distribution of zeros

- $\zeta(s) \neq 0$ for $\mathfrak{Re}(s) = 1$: $\pi(x)$, $\pi_{a,q}(x)$.
- GRH: error terms.
- GSH: Chebyshev's bias.
- Analytic rank, adjacent spacings: *h*(*D*).

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Explicit Formula (Contour Integration)

$$-\frac{\zeta'(s)}{\zeta(s)} = -\frac{\mathsf{d}}{\mathsf{d}s}\log\zeta(s) = -\frac{\mathsf{d}}{\mathsf{d}s}\log\prod_{p}\left(1-p^{-s}\right)^{-1}$$

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Explicit Formula (Contour Integration)

$$\frac{\zeta'(s)}{\zeta(s)} = -\frac{d}{ds} \log \zeta(s) = -\frac{d}{ds} \log \prod_{p} (1-p^{-s})^{-1}$$
$$= \frac{d}{ds} \sum_{p} \log (1-p^{-s})$$
$$= \sum_{p} \frac{\log p \cdot p^{-s}}{1-p^{-s}} = \sum_{p} \frac{\log p}{p^{s}} + \operatorname{Good}(s).$$

Explicit Formula (Contour Integration)

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Contour Integration:

$$\int -\frac{\zeta'(s)}{\zeta(s)} \frac{x^s}{s} ds \quad \text{vs} \quad \sum_p \log p \int \left(\frac{x}{p}\right)^s \frac{ds}{s}$$

Ref

Explicit Formula (Contour Integration)

$$-\frac{\zeta'(s)}{\zeta(s)} = -\frac{d}{ds}\log\zeta(s) = -\frac{d}{ds}\log\prod_{p}\left(1-p^{-s}\right)^{-1}$$
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Contour Integration:

$$\int - \frac{\zeta'(s)}{\zeta(s)} \phi(s) ds$$
 vs $\sum_p \log p \int \phi(s) p^{-s} ds.$

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Explicit Formula (Contour Integration)

$$-\frac{\zeta'(s)}{\zeta(s)} = -\frac{d}{ds}\log\zeta(s) = -\frac{d}{ds}\log\prod_{p}\left(1-p^{-s}\right)^{-1}$$
$$= \frac{d}{ds}\sum_{p}\log\left(1-p^{-s}\right)$$
$$= \sum_{p}\frac{\log p \cdot p^{-s}}{1-p^{-s}} = \sum_{p}\frac{\log p}{p^{s}} + \operatorname{Good}(s).$$

Contour Integration (see Fourier Transform arising):

$$\int -rac{\zeta'(s)}{\zeta(s)} \, \phi(s) ds$$
 vs $\sum_p \log p \int \phi(s) e^{-\sigma \log p} e^{-it \log p} ds.$

Knowledge of zeros gives info on coefficients.

Modeling lowest zero of $L_{E_{11}}(s, \chi_d)$ with 0 < d < 400,000



of SO(2N) with $N_{\rm eff}$ (solid), standard N_0 (dashed).

Modeling lowest zero of $L_{E_{11}}(s, \chi_d)$ with 0 < d < 400,000



Lowest zero for $L_{E_{11}}(s, \chi_d)$ (bar chart), lowest eigenvalue of SO(2N) with $N_0 = 12$ (solid) with discretisation and with standard $N_0 = 12.26$ (dashed) without discretisation.

Modeling lowest zero of $L_{E_{11}}(s, \chi_d)$ with 0 < d < 400,000



Lowest zero of $L_{E_{11}}(s, \chi_d)$ (bar chart), lowest eigenvalue of SO(2N) effective *N* of N_{eff} = 2 (solid) w' discretisation and effective *N* of N_{eff} = 2.32 (dashed) w/o discretisation.

Correspondences

Similarities between *L*-Functions and Nuclei:

Zeros \longleftrightarrow Energy Levels

Schwartz test function \longrightarrow Neutron

Support of test function \leftrightarrow Neutron Energy.



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Research Experiences for Undergraduates

Williams College SMALL REU: Deadline Wednesday February 3rd: 5pm US Eastern: https: //math.williams.edu/small/

Polymath Jr: Deadline April 1st: 5pm US Eastern: https://geometrynyc. wixsite.com/polymathreu

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